

# Nuclear Physics: BSM

*William Detmold*



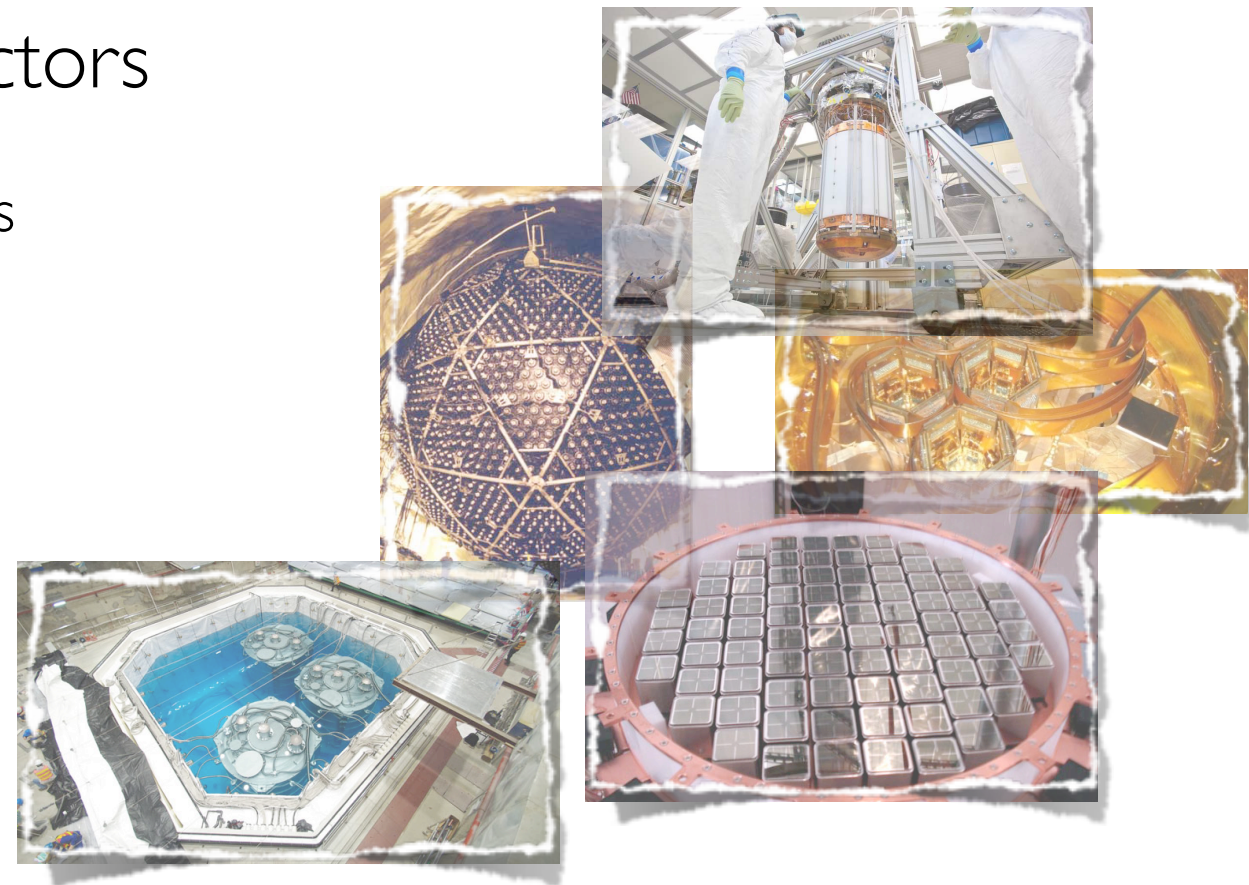
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Institute of  
Technology

# Outline

- Role of NP in search for physics BSM and other rare processes
- Uncertainties in nuclear effects: (un)known (un)knowns
  - Examples:  $0\nu\beta\beta$  decay, GT transitions, eA DIS
- Nuclear physics from the SM?
  - Problems and progress
  - Scalar-isoscalar current nuclear matrix elements

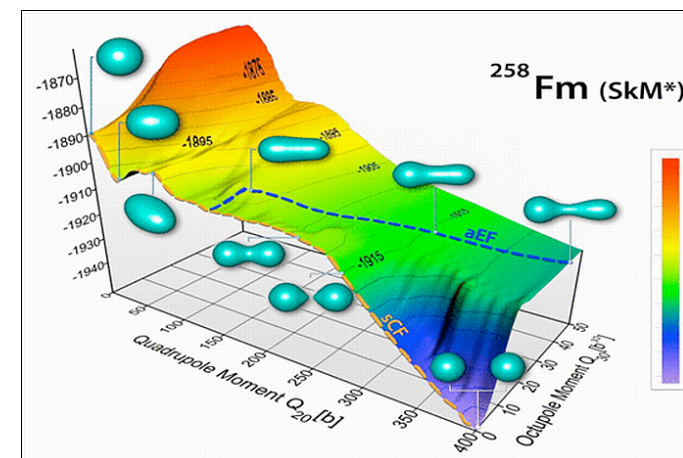
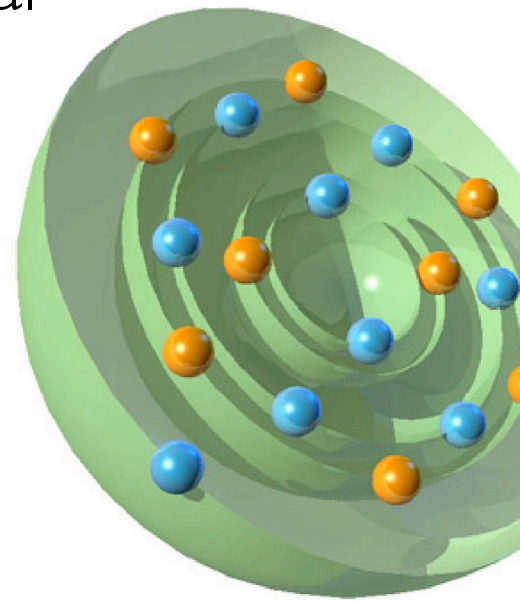
# *Searches for very rare interactions*

- Weakness of weak or dark matter interactions with ordinary matter  $\Rightarrow$  big detectors
- Piles (buckets) of heavy stuff instead of protons
- Specific nuclear transitions sensitive to fundamental symmetries: nuclear environment enhances effect
- Range of nuclei used in detectors
  - Neutrino DIS: NuTeV - steel sheets
  - Solar Neutrinos @ SNO: D<sub>2</sub>O
  - Dark matter: Na, Si, Ge, Xe, ...
  - EDM searches: Rn, Hg, Xe, ..
  - $0\nu\beta\beta$  decay: Ge, Ca, Xe, ...



# How predictive is nuclear physics?

- NP critical to experimental investigations of these interactions
- Phenomenological: poorly understood from a theoretical perspective
- Quantifying uncertainties on predictions is difficult
- Ideal experiment: simple nuclei (single isotope) where theory is best understood
- Study dependence on target
- Reality: single target chosen for a host of reasons
- 😊 New physics *discovery* can be robust despite NP
- 😞 Differentiation of new physics no so easy

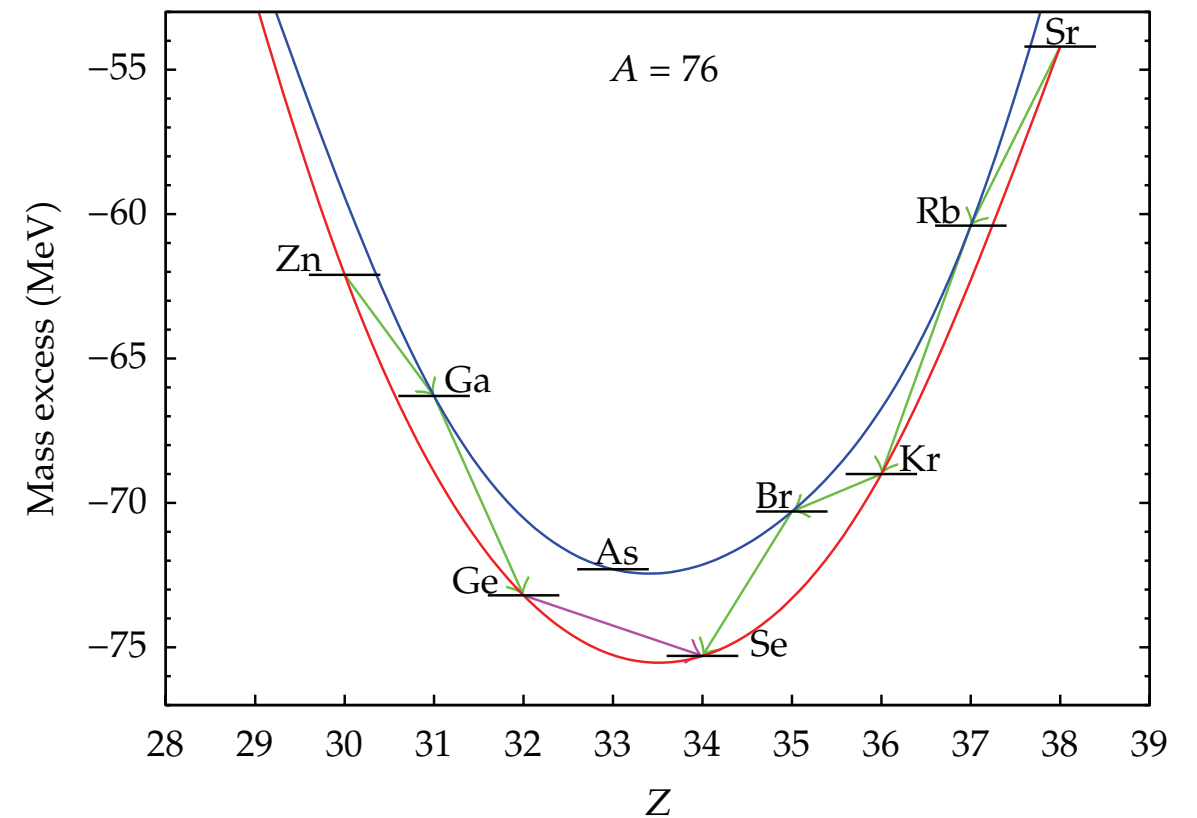
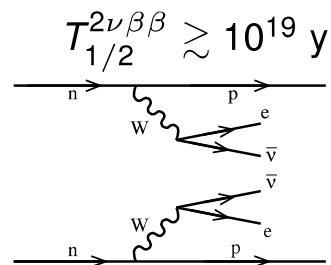




$0\nu\beta\beta$  decay

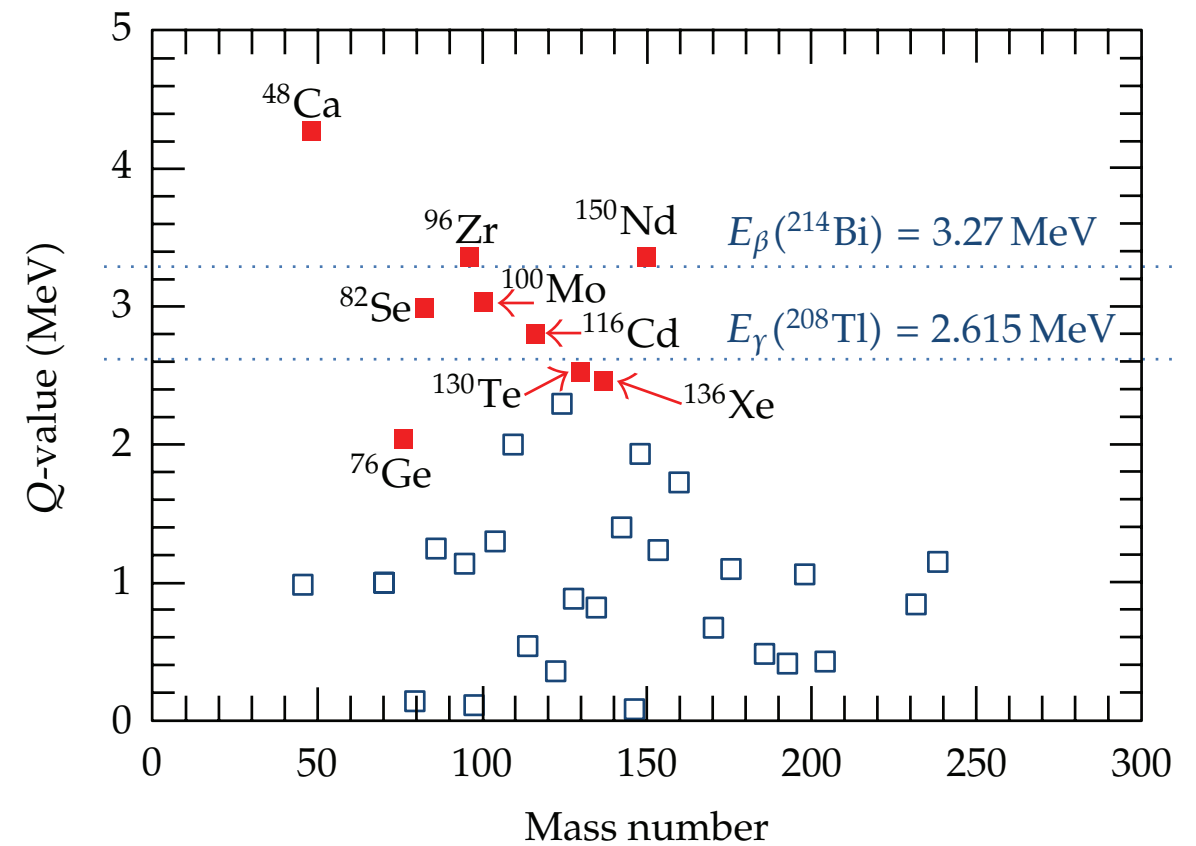
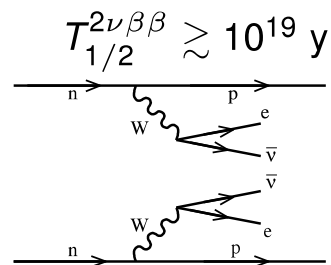
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- Certain nuclei allow observable  $\beta\beta$  decay



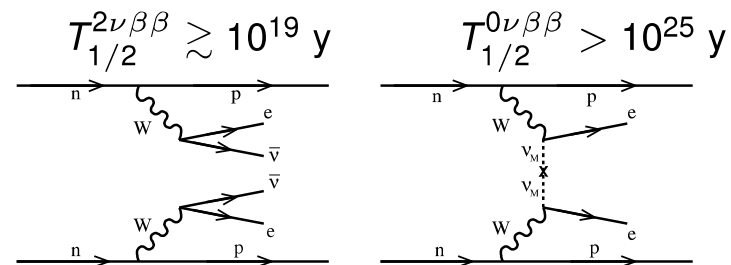
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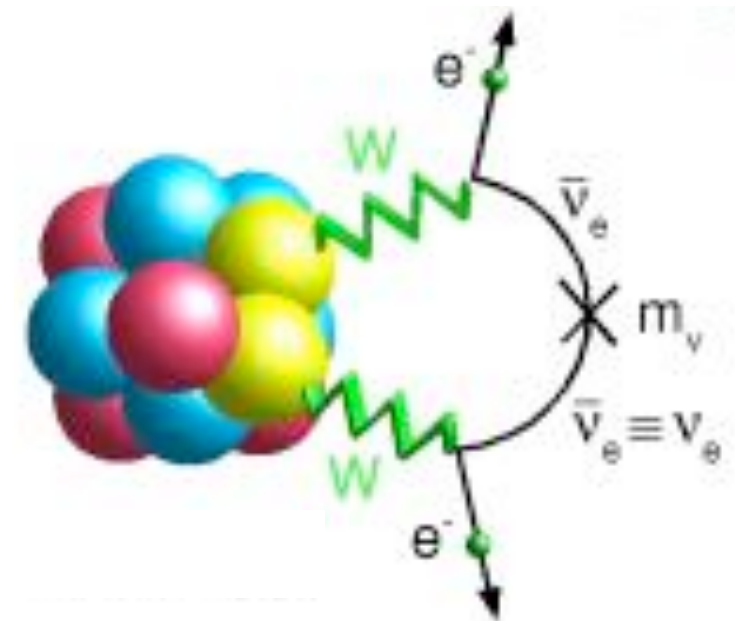
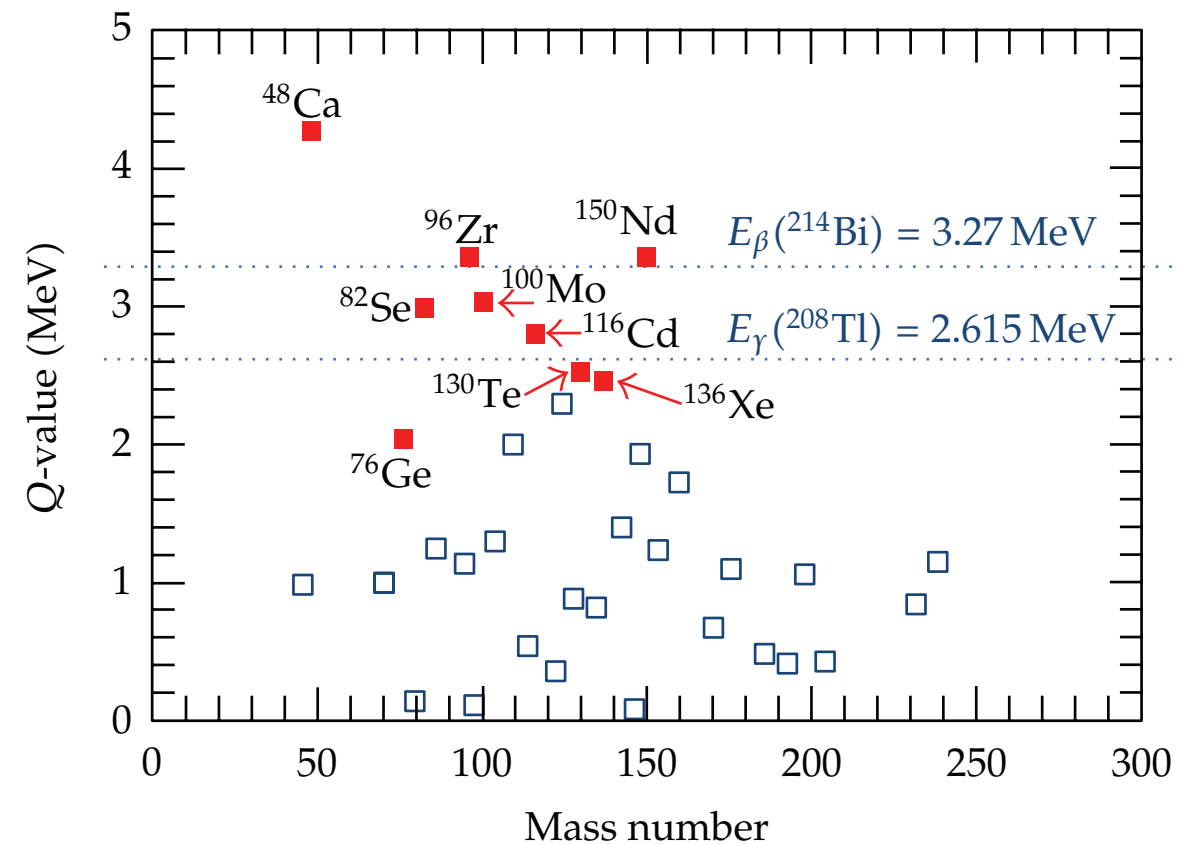


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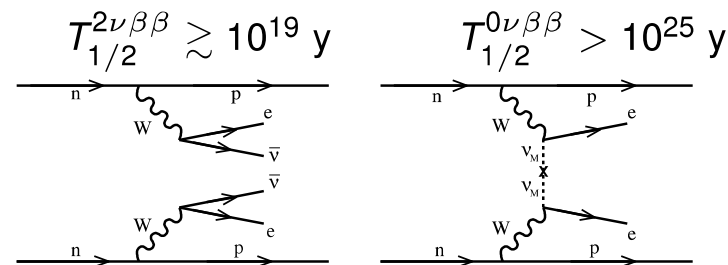
- If neutrinos are massive Majorana fermions  $0\nu\beta\beta$  decay is possible





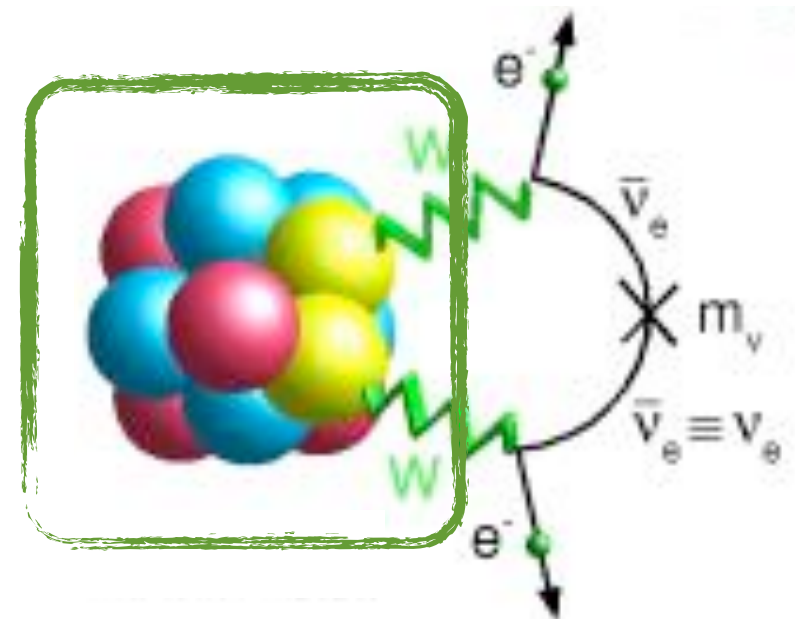
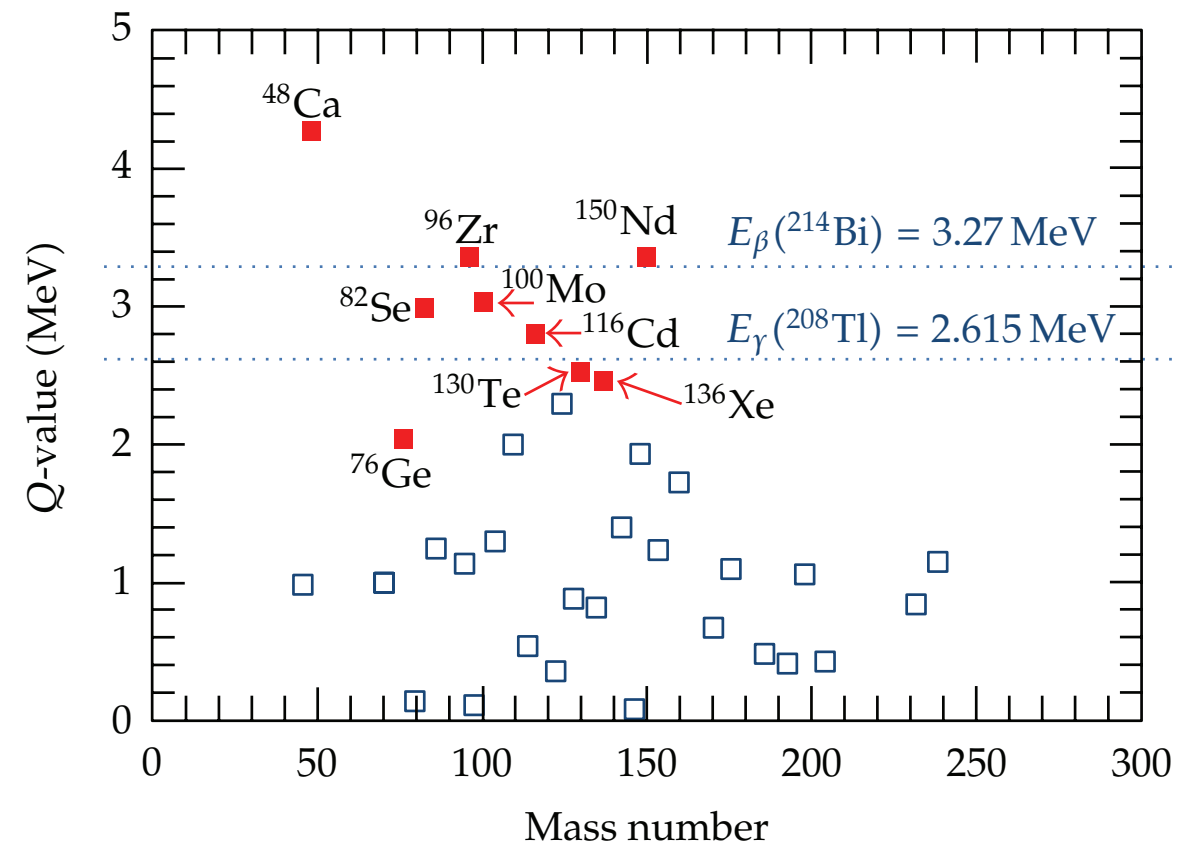
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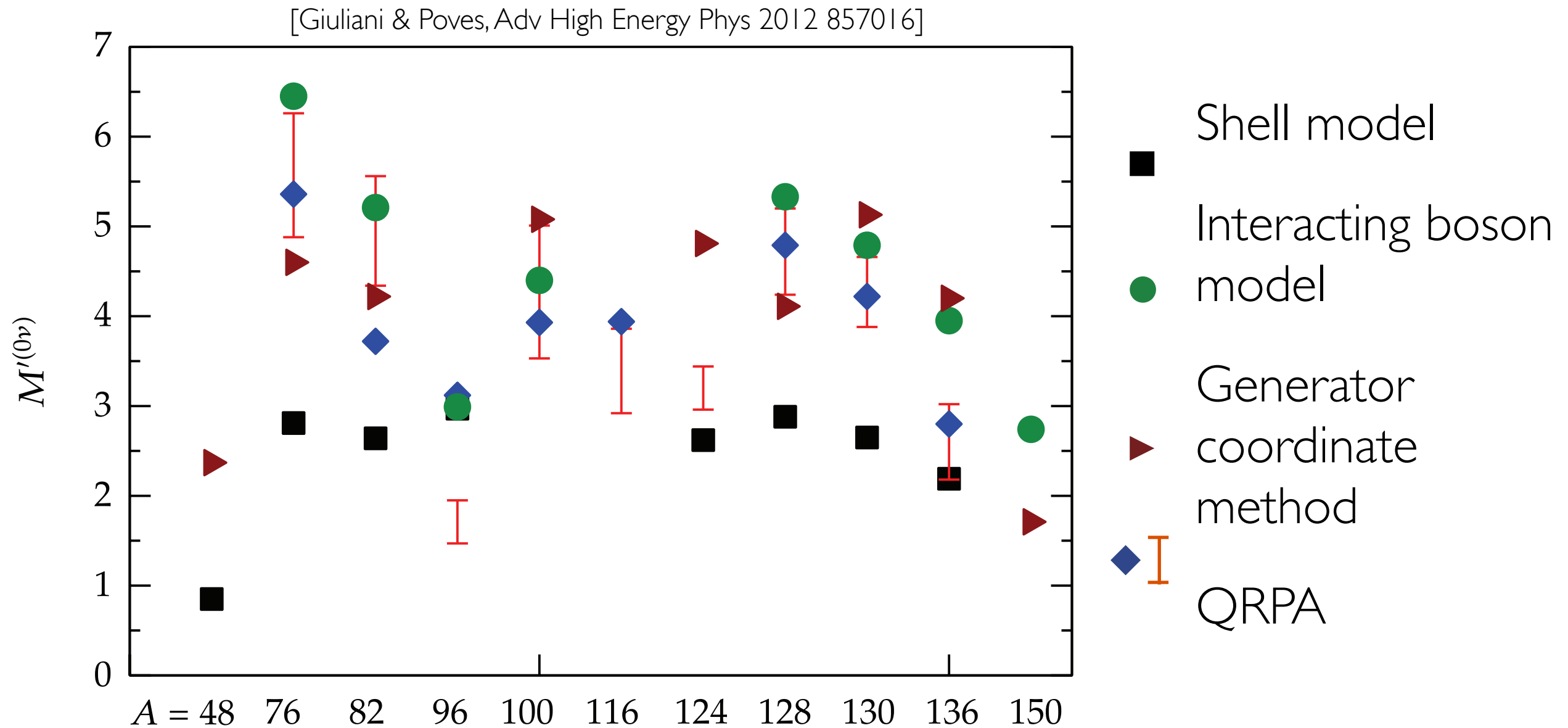


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- Half-life depends critically on the nuclear matrix elements of two weak currents
- $$\left(T_{1/2}^{0\nu\beta\beta}(0^+ \rightarrow 0^+)\right)^{-1} = G_{01} \left|M^{0\nu\beta\beta}\right|^2 \left(\frac{\langle m_\nu \rangle}{m_e}\right)^2$$



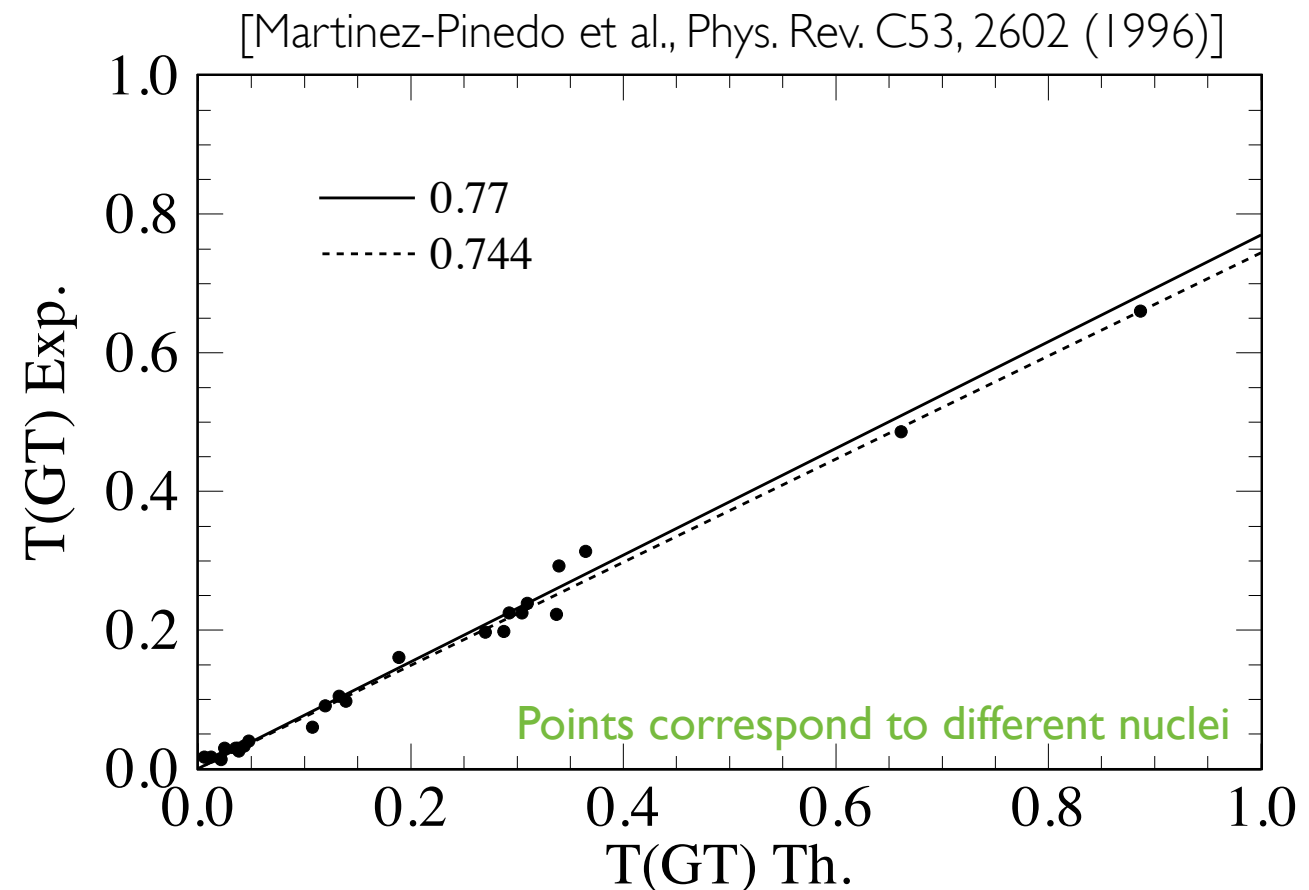
# $0\nu\beta\beta$ decay nuclear matrix elements



Is the spread of results representative of the true uncertainty?

# Gamow-Teller: axial charge in nuclei

- Gamow-Teller transitions in nuclei are a stark example of problems
- Well measured
- Best nuclear structure calculations are systematically off by 20–30%
- Large range of nuclei ( $30 < A < 60$ ) where spectrum is well described
- QRPA, shell-model, ...
- Correct for it by “quenching” axial charge in nuclei ...

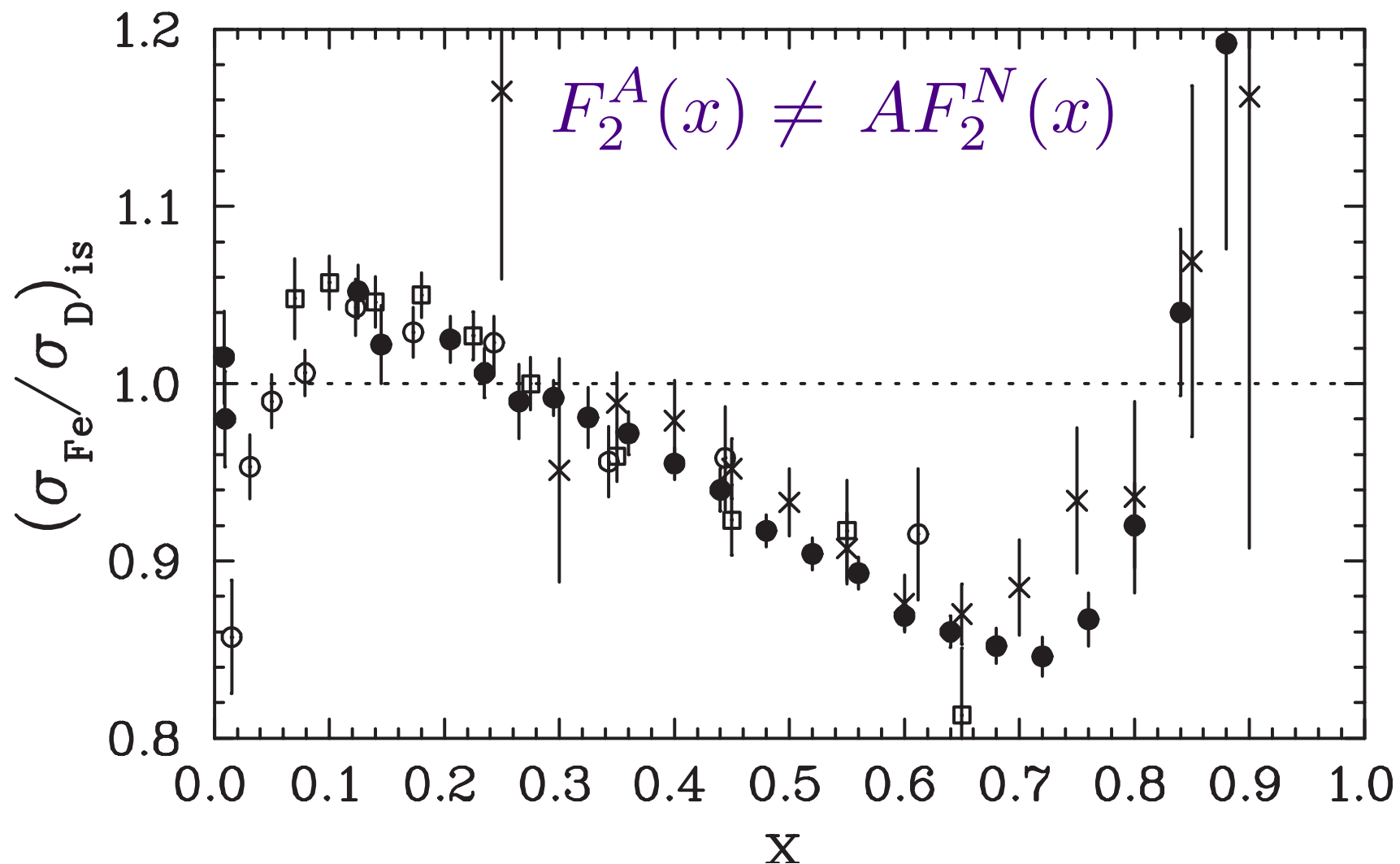


$$T(GT) \sim \sqrt{\sum_f \langle \boldsymbol{\sigma} \cdot \boldsymbol{\tau} \rangle_{i \rightarrow f}}$$

$$\langle \boldsymbol{\sigma} \boldsymbol{\tau} \rangle = \frac{\langle f || \sum_k \boldsymbol{\sigma}^k \boldsymbol{t}_{\pm}^k || i \rangle}{\sqrt{2J_i + 1}}$$

# The EMC effect

- 1983: DIS on Fe target [EMC]

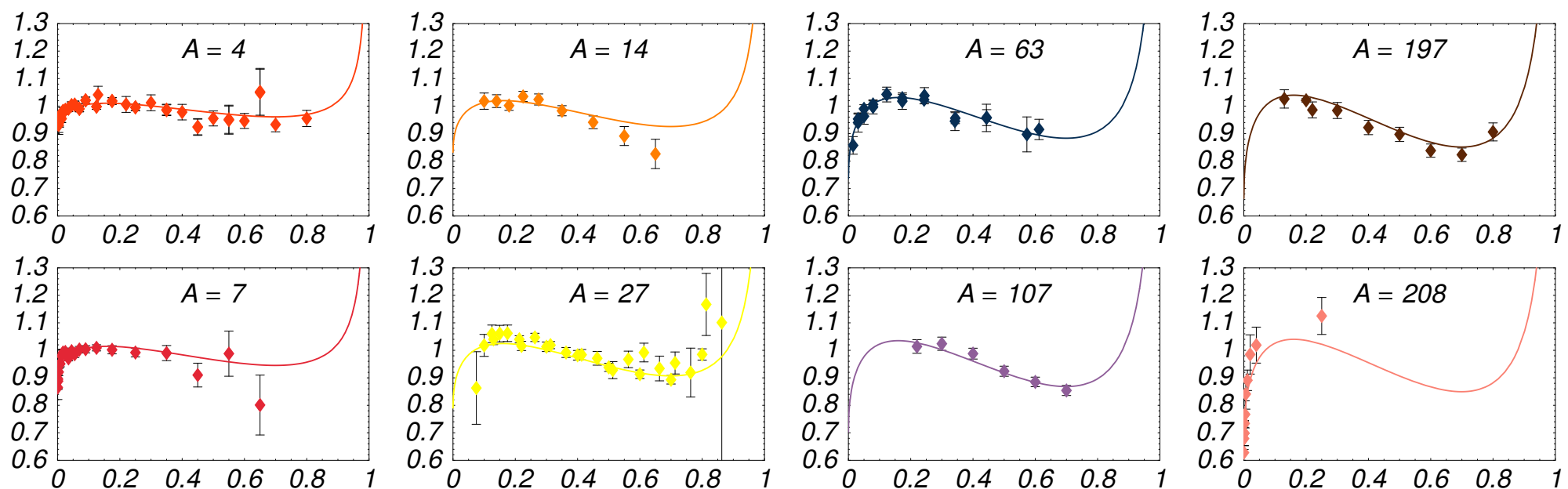


- Proton structure modified in a nuclear environment



# Dependence on $A$

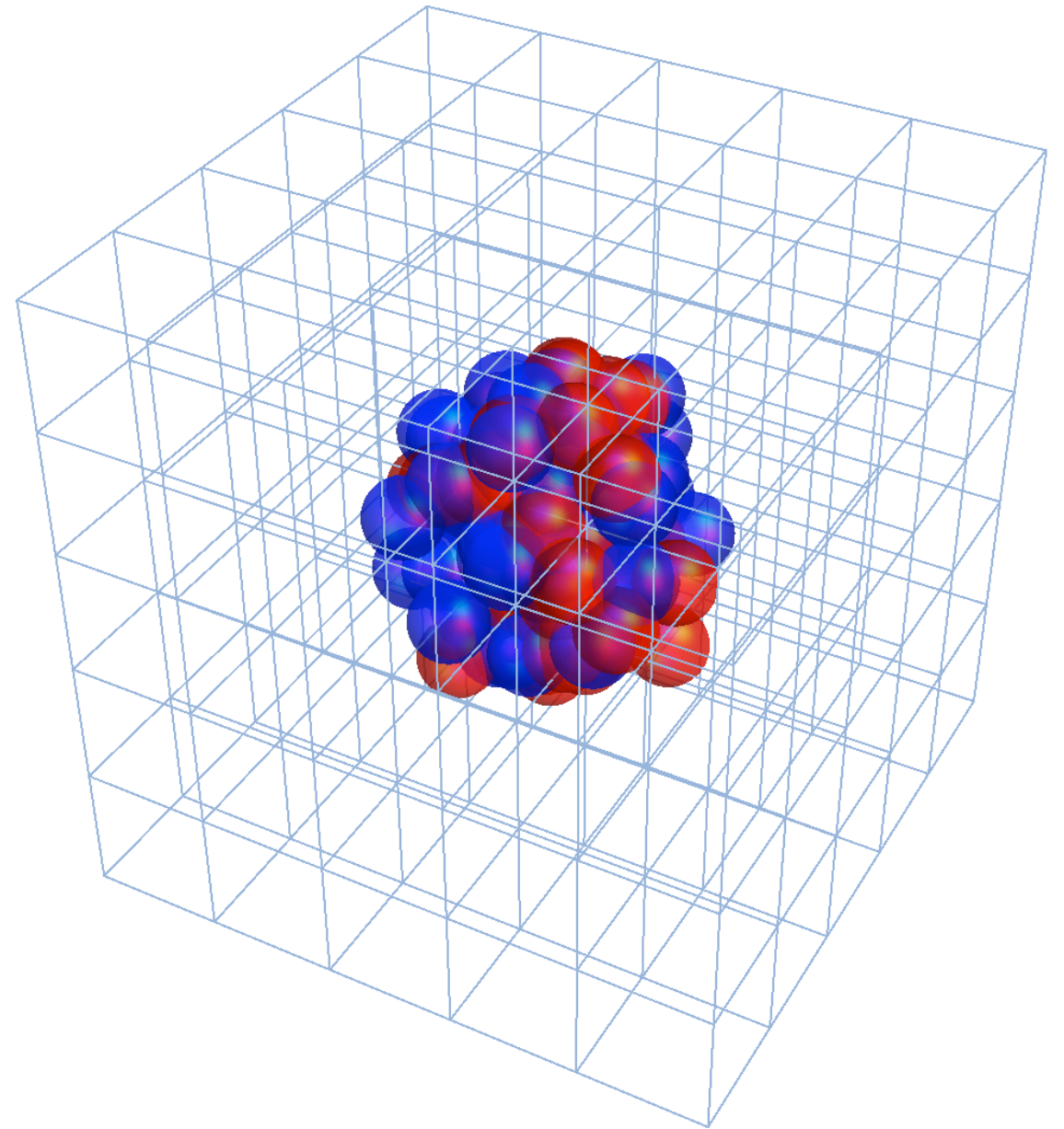
- Large nuclear effects  $\sim 30\%$
- Over last 30 years: studies of target dependence



- No convincing microscopic understanding of its origin (EMC = Every Model is Cool)
- Little predictive power

# *LQCD to the rescue?*

- Nuclear physics is Standard Model physics
- ... so calculate ab initio!



*Nuclei: an (exponentially hard)<sup>2</sup> problem*

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 $\langle 0 | T q_1(t) \dots q_{624}(t) \bar{q}_1(0) \dots \bar{q}_{624}(0) | 0 \rangle$



# Nuclei: an (exponentially hard)<sup>2</sup> problem

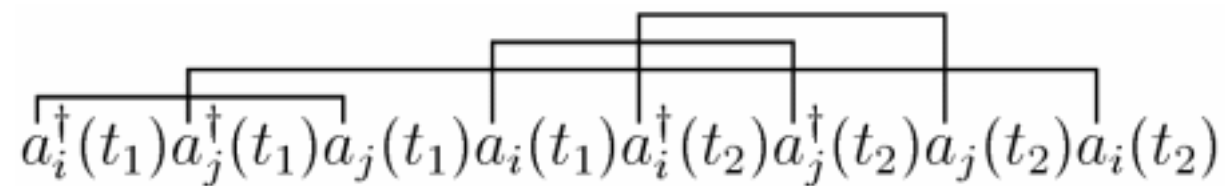
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$$\langle 0 | T q_1(t) \dots q_{624}(t) \bar{q}_1(0) \dots \bar{q}_{624}(0) | 0 \rangle$$
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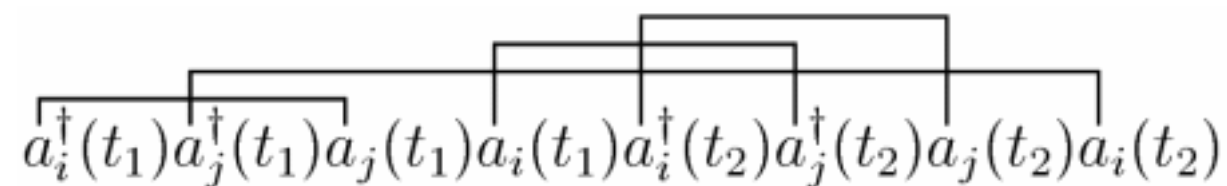
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- Complexity: number of contractions =  $(A+Z)!(2A-Z)!$



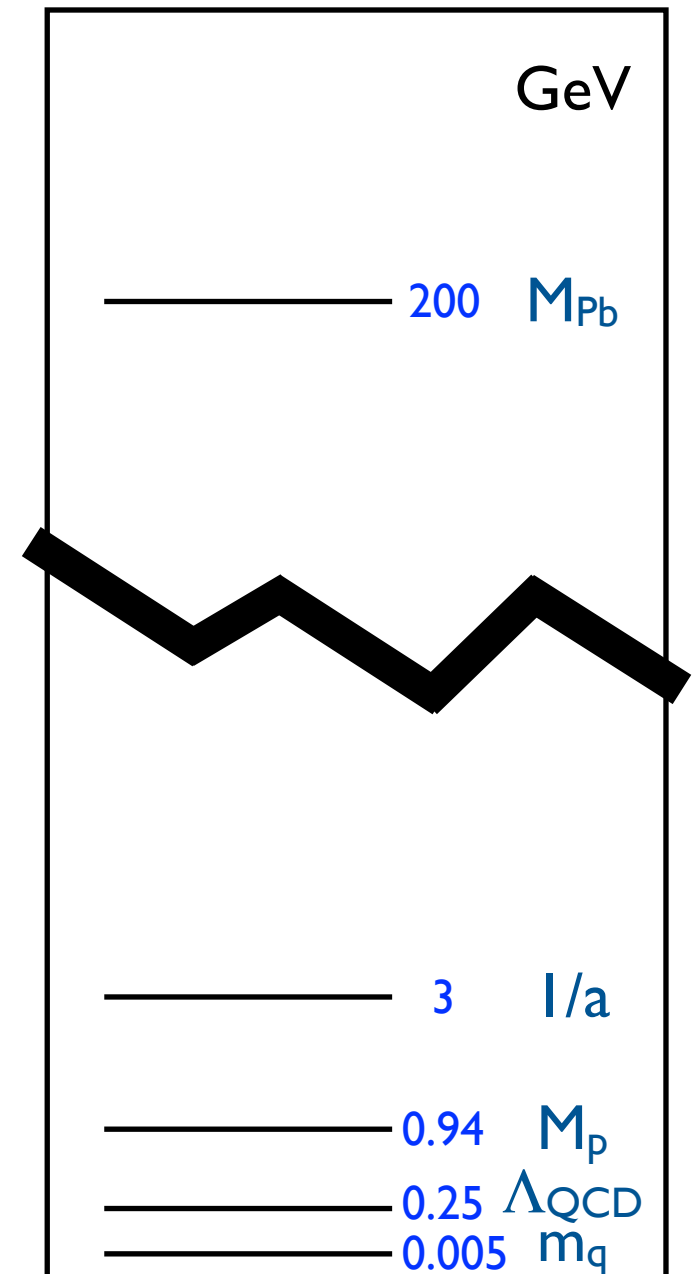
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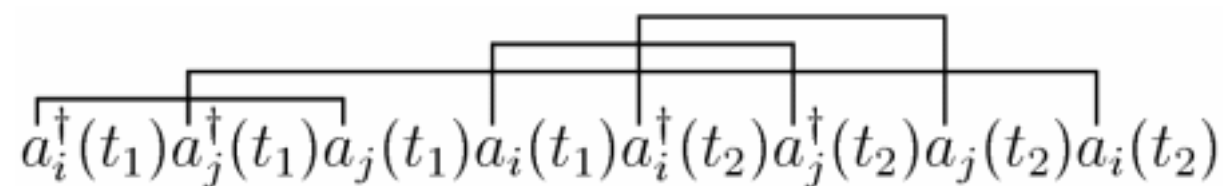
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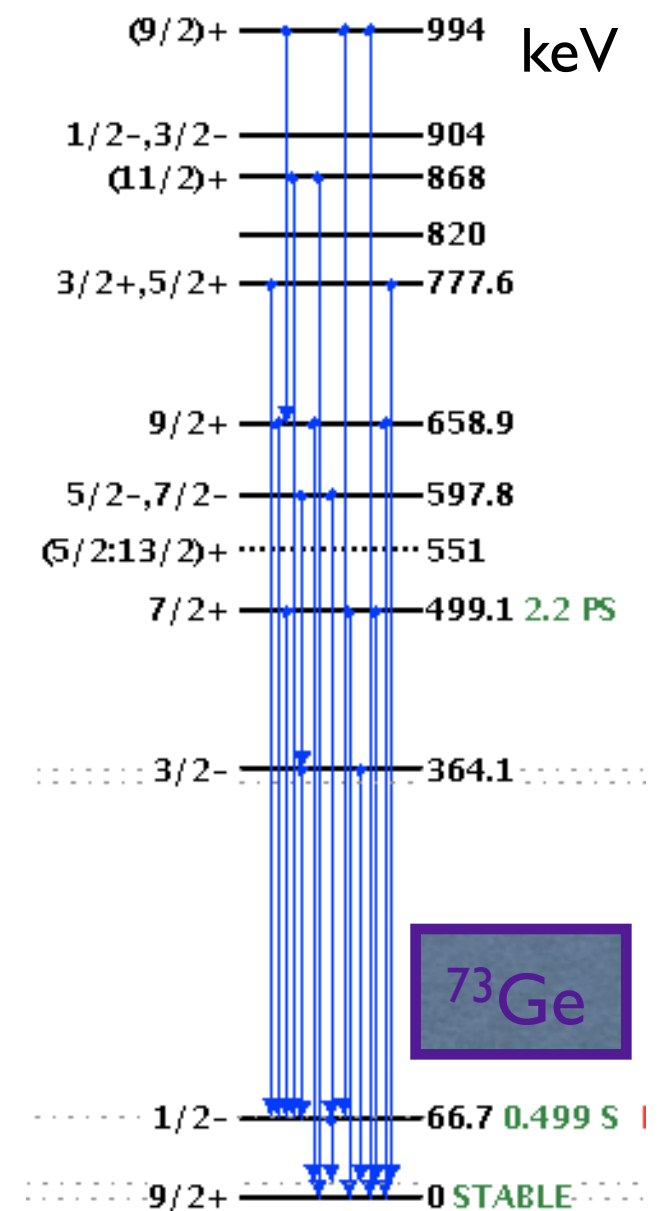
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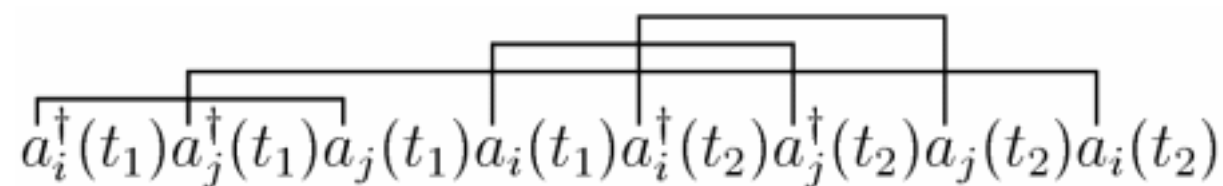
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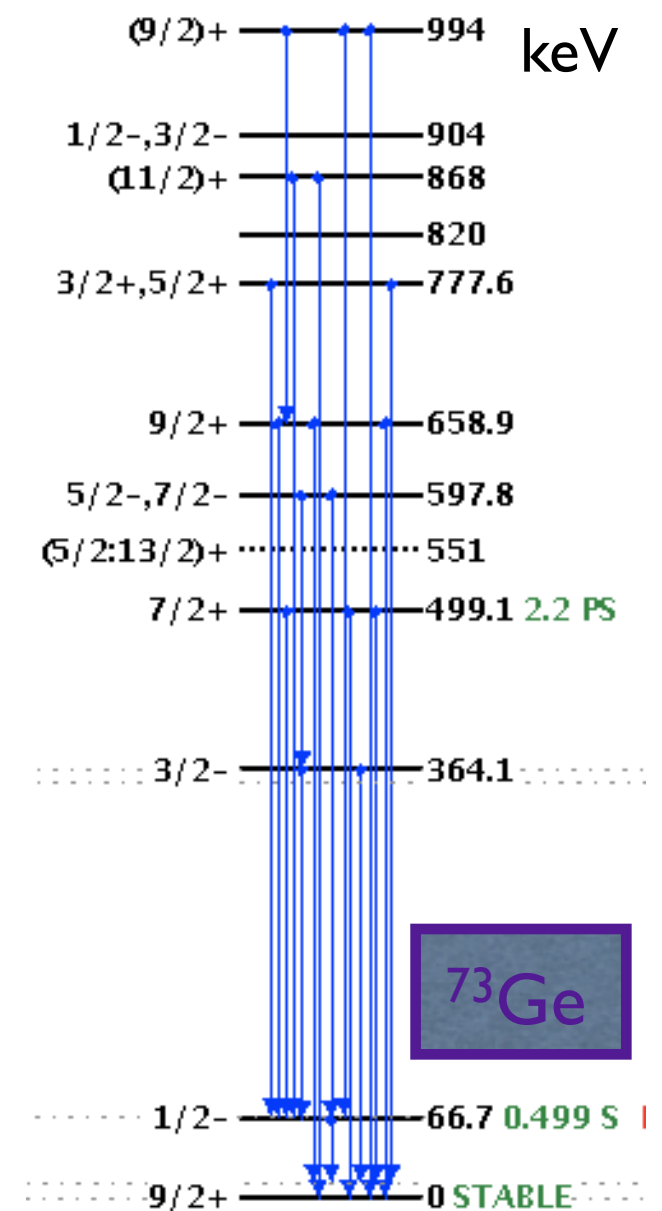
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- Small energy splittings
- Importance sampling: statistical noise exponentially increases with  $A$



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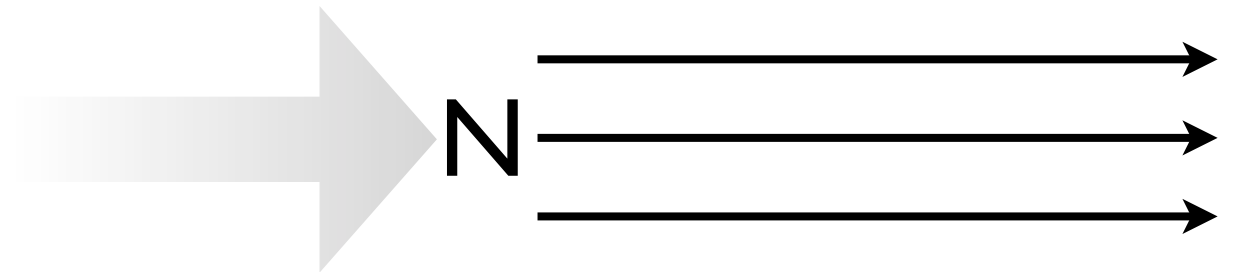
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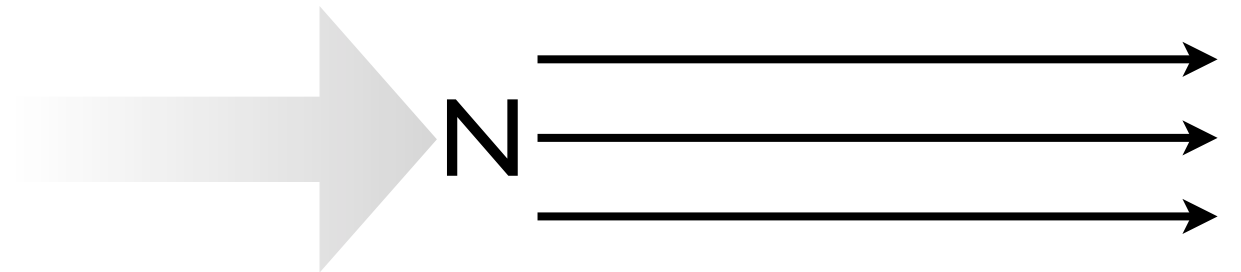


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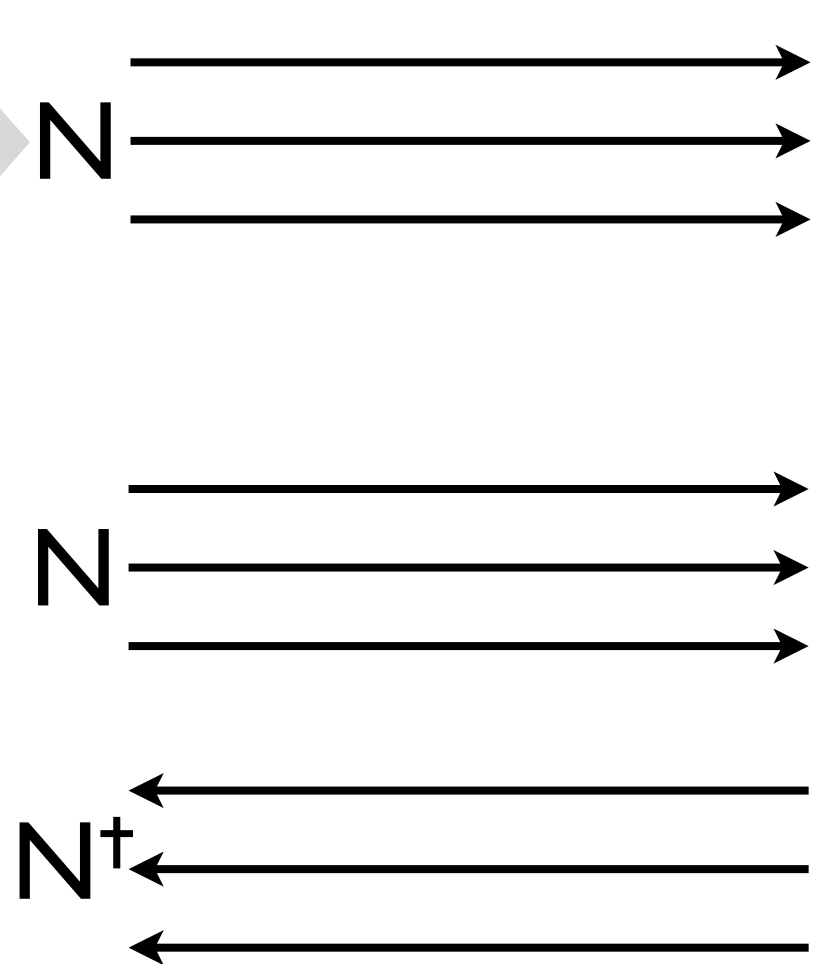
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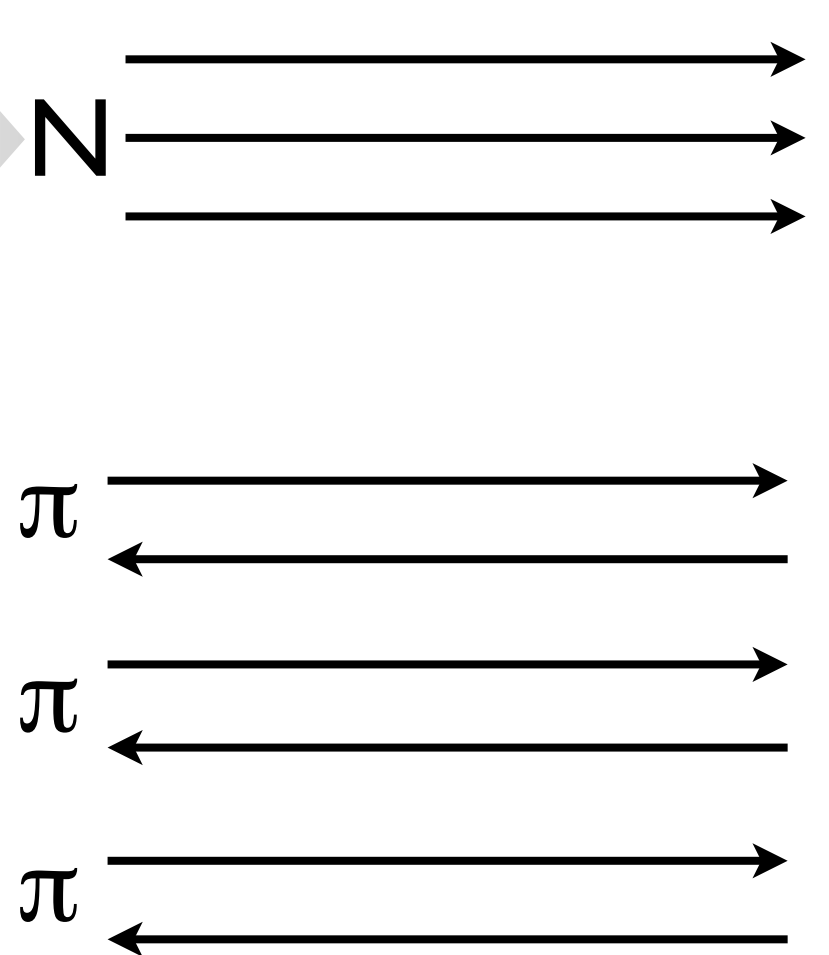
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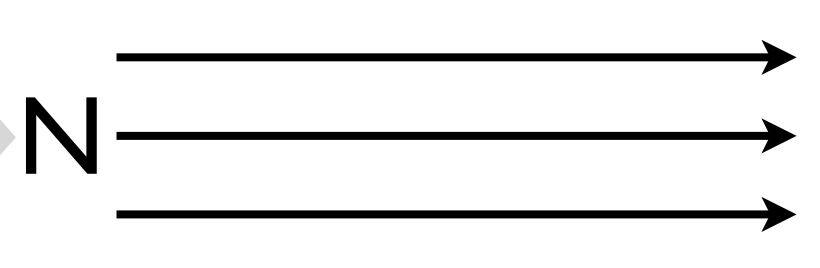
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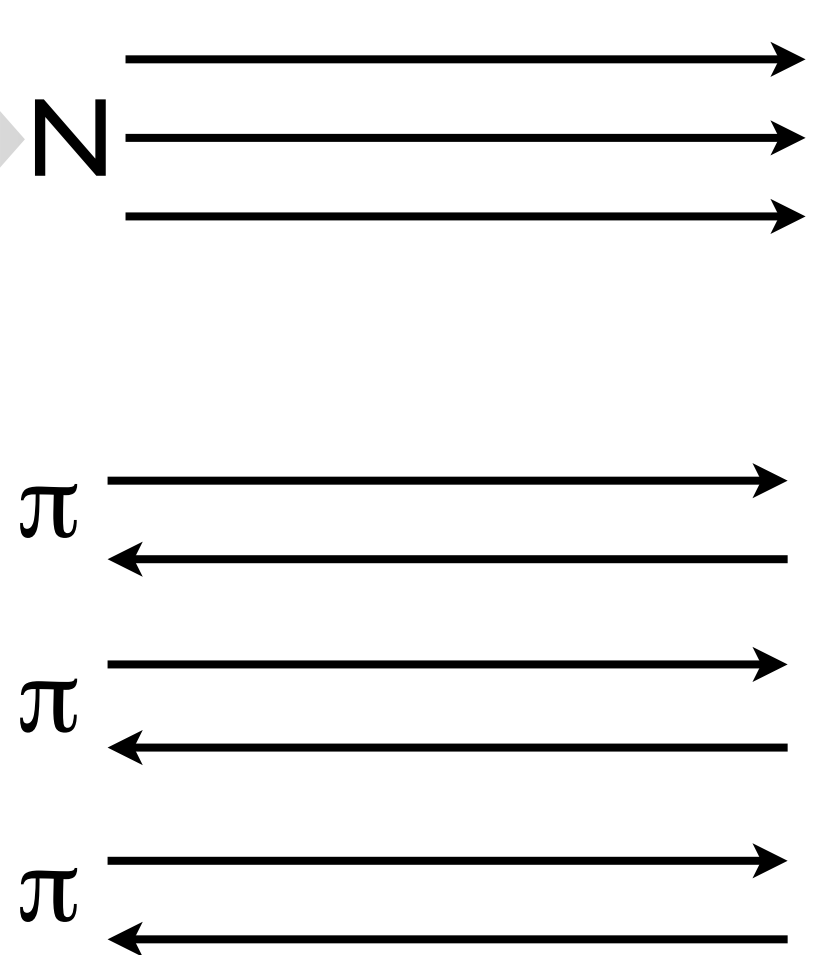
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$$\frac{\text{signal}}{\text{noise}} \sim \exp[-(M_N - 3/2 m_\pi)t]$$



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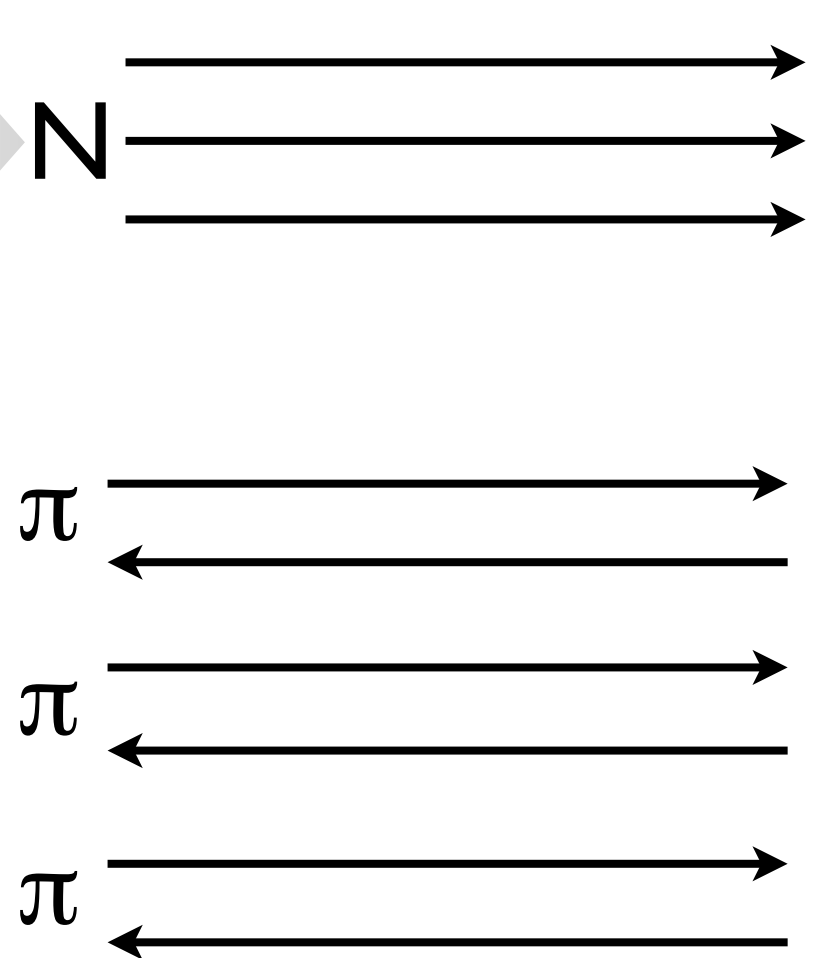
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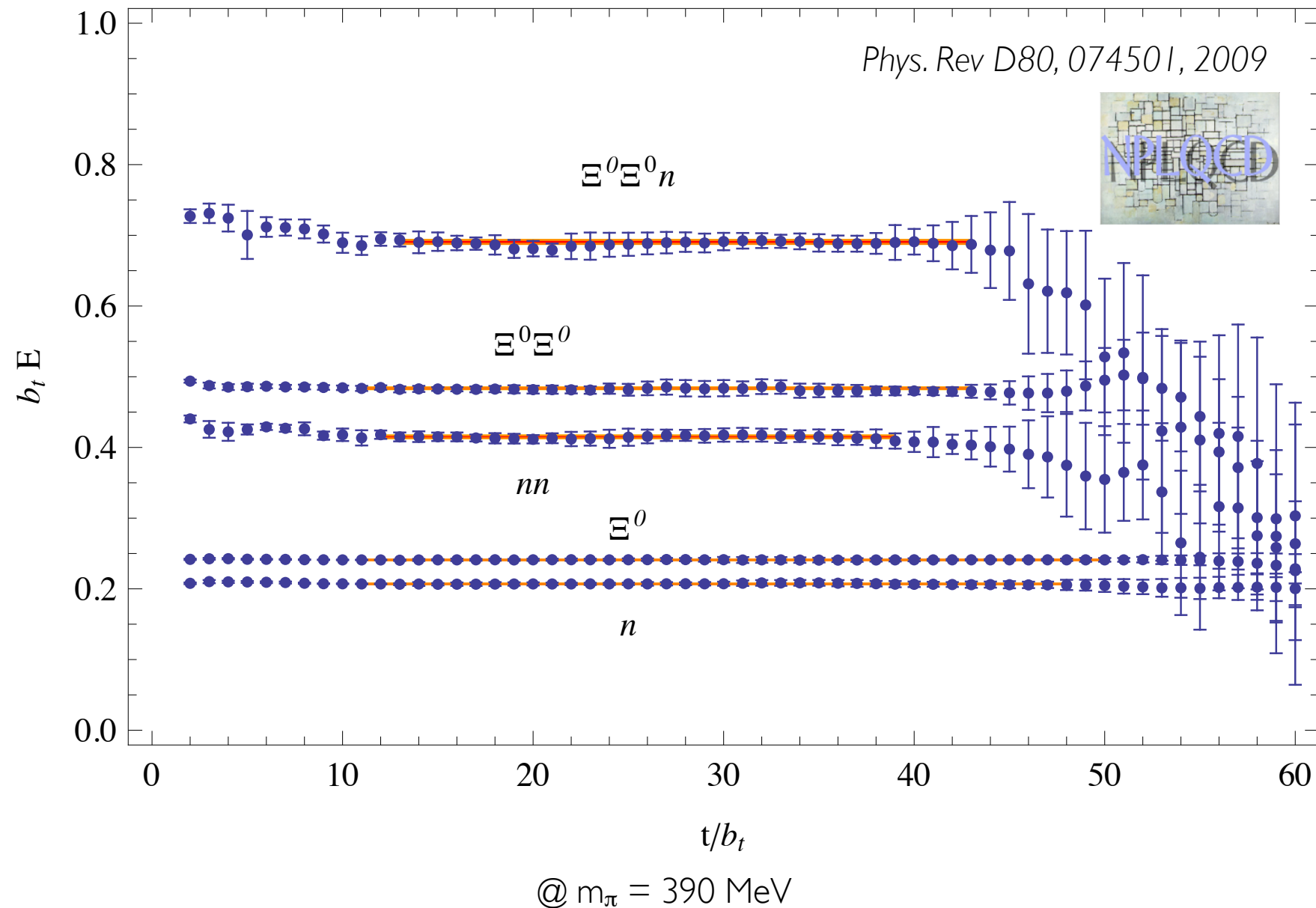
- For nucleus A:

$$\frac{\text{signal}}{\text{noise}} \sim \exp[-\textcolor{red}{A}(M_N - 3/2 m_\pi)t]$$



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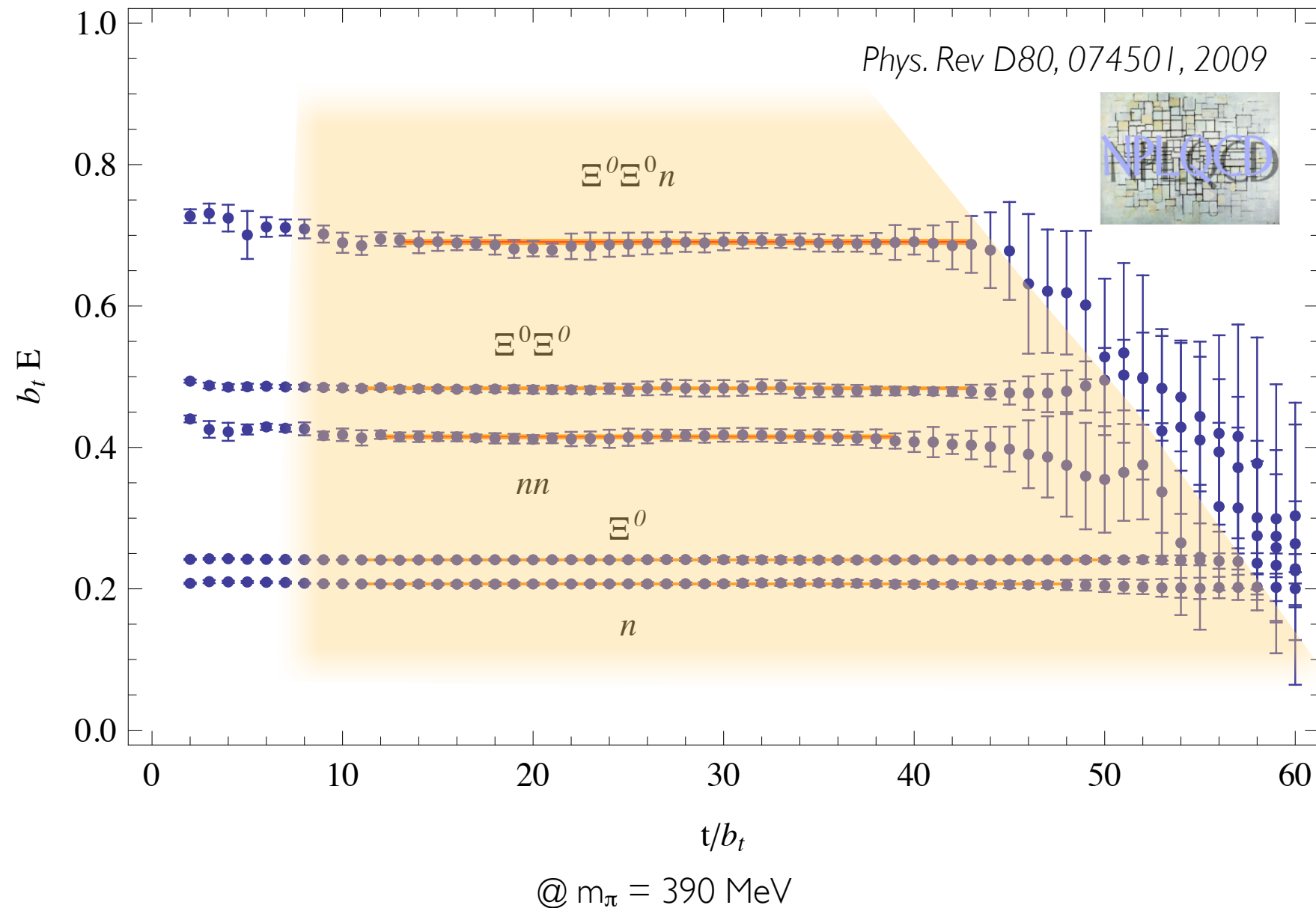
High statistics study using anisotropic lattices (fine temporal resolution)





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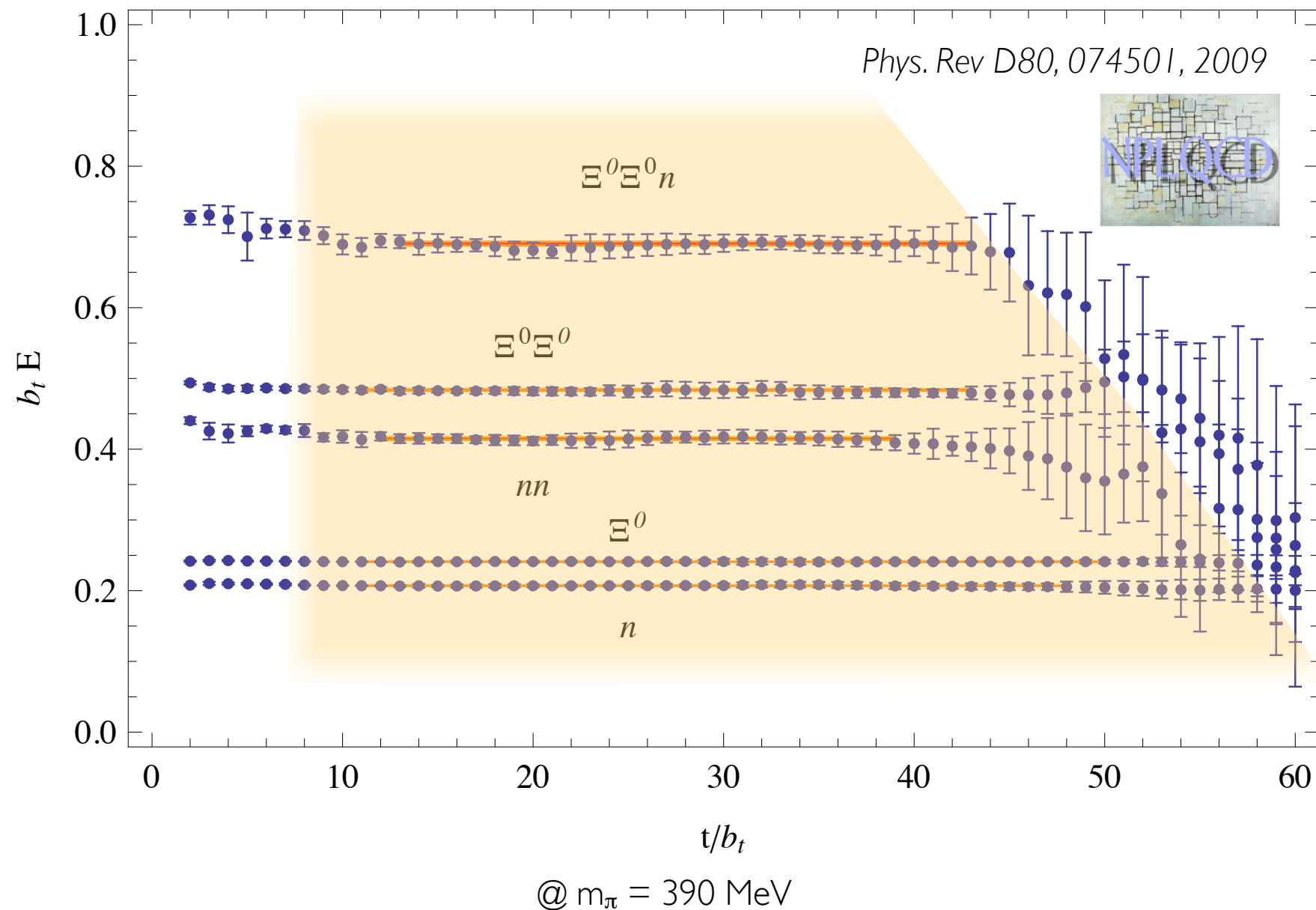
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Golden window of time-slices where signal/noise const

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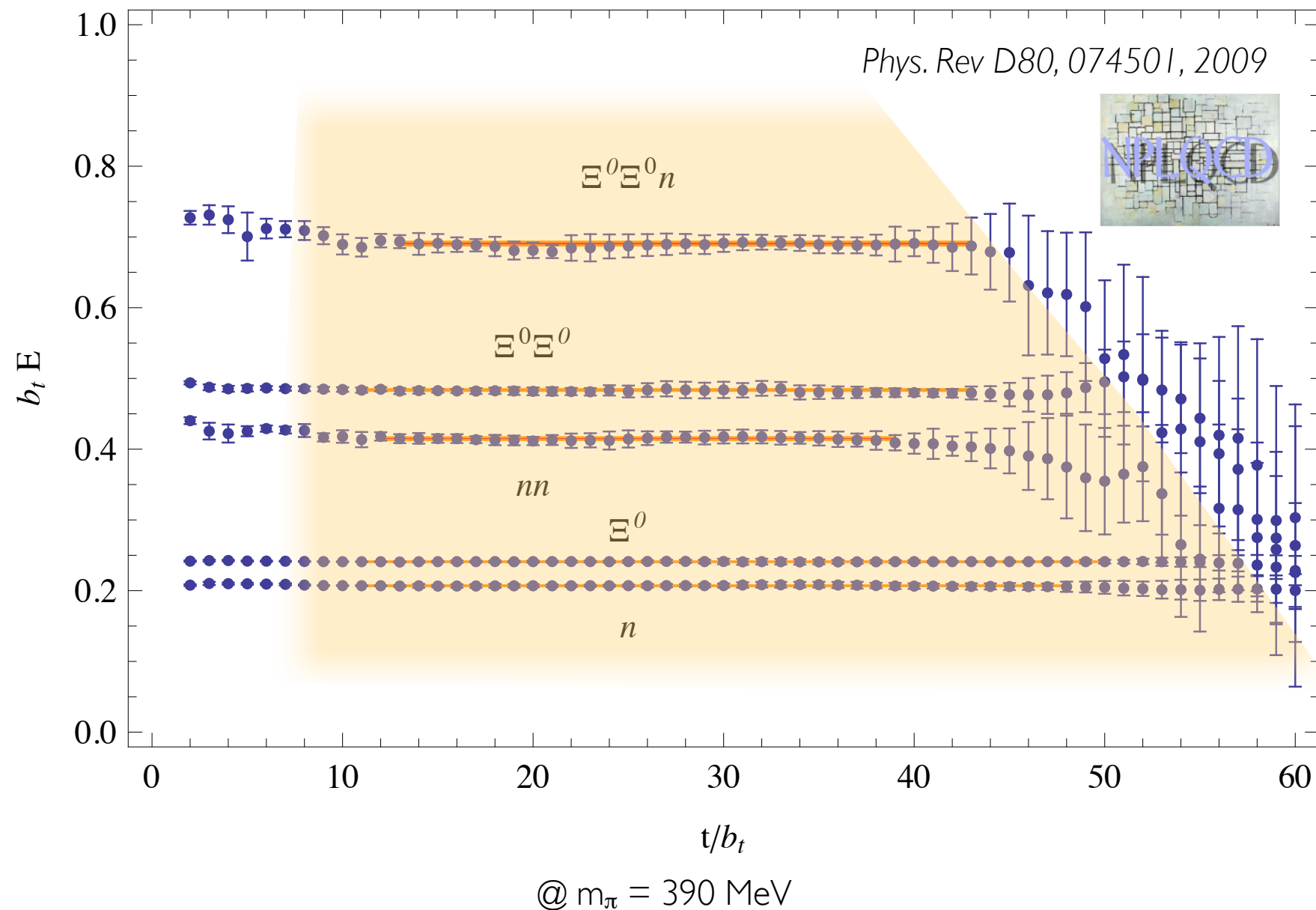
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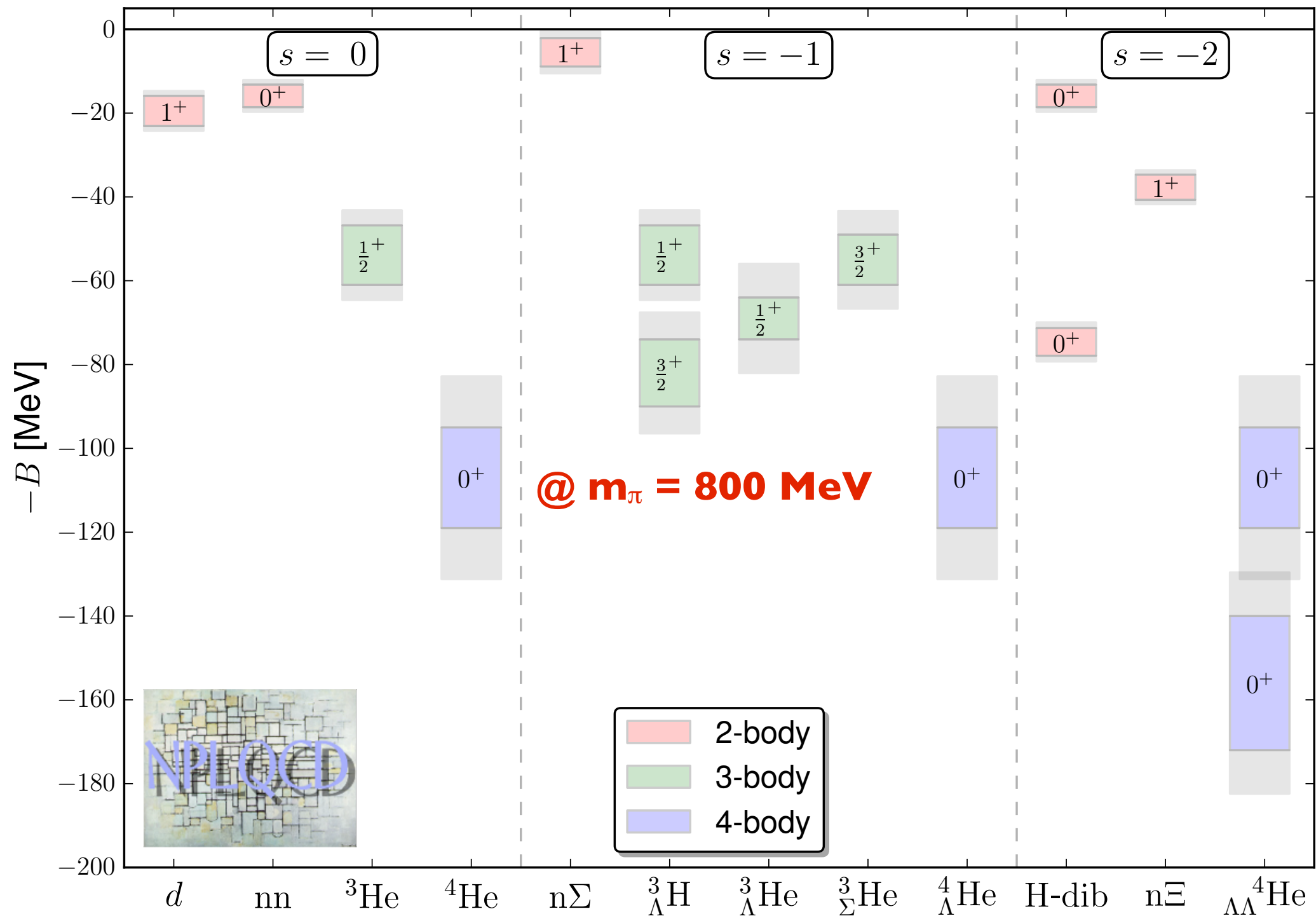
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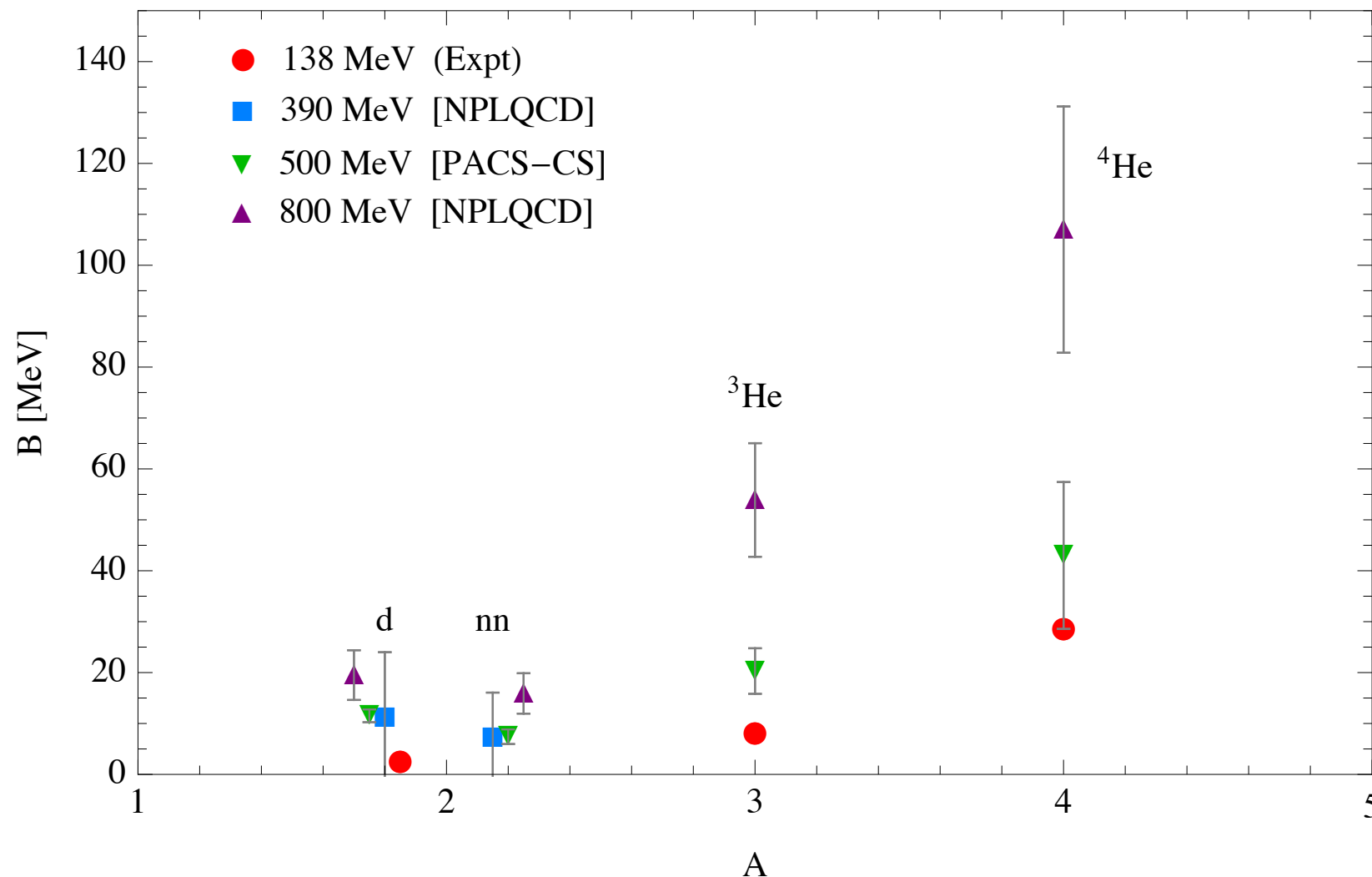
Golden window of time-slices where signal/noise const

Interpolator choice can be used to suppress noise

# Nuclei ( $A=2,3,4$ )



# QCD Periodic Table

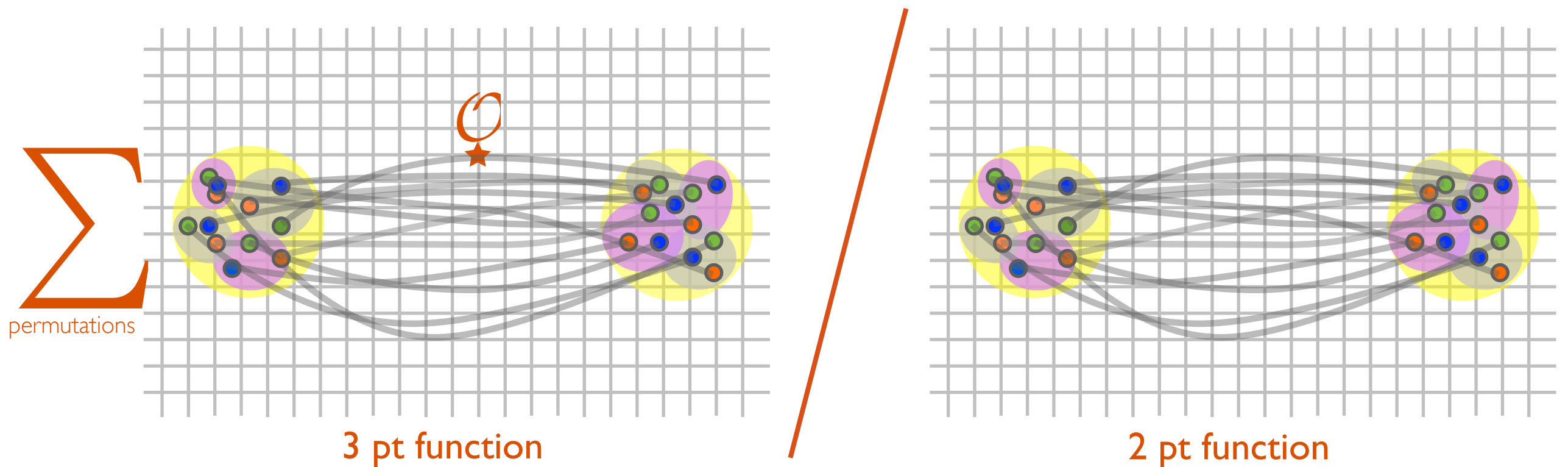


- Light nuclei ( $A < 6$ ) will be feasible at the physical quark mass in the near future
- Interesting progress with larger nuclei ( $A=4, 8, 12, \dots$ ) but still a major challenge [WD, Orginos Phys.Rev. D87 (2013) 114512]



# Nuclear matrix elements

- Calculations of matrix elements of currents in light nuclei just beginning
- For deeply bound nuclei, use the same techniques as for single hadron matrix elements

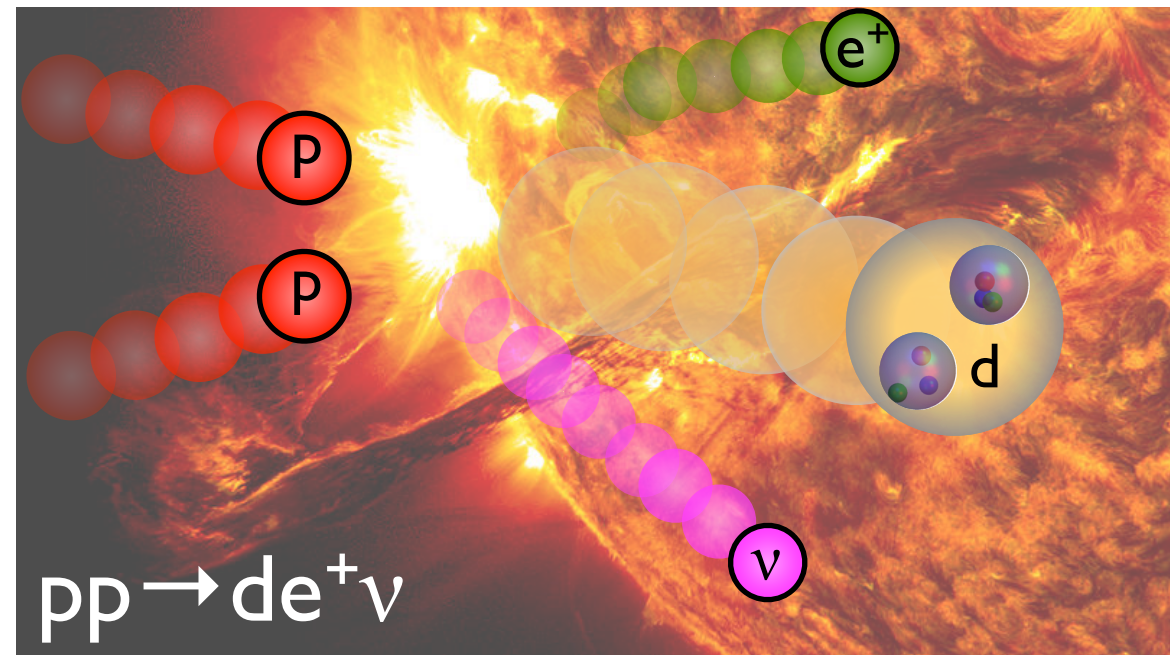


- For near threshold states, need to be careful with volume effects



# Nuclear matrix elements

1. Axial coupling to NN system
  - pp fusion: “Calibrate the sun”
  - Muon capture: MuSun @ PSI
  - $d\nu \rightarrow nne^+ : \text{SNO}$



2. Twist-2 operators: eg EMC effect

$$\langle N, Z | \bar{q} \gamma_{\{\mu_0} D_{\mu_1} \dots D_{\mu_n\}} q | N, Z \rangle$$

- Velocity dependent DM interactions
- Proof of principle (moments of pion PDF in pion gas) [WD, HW Lin 1112.5682]

# Nuclear sigma terms

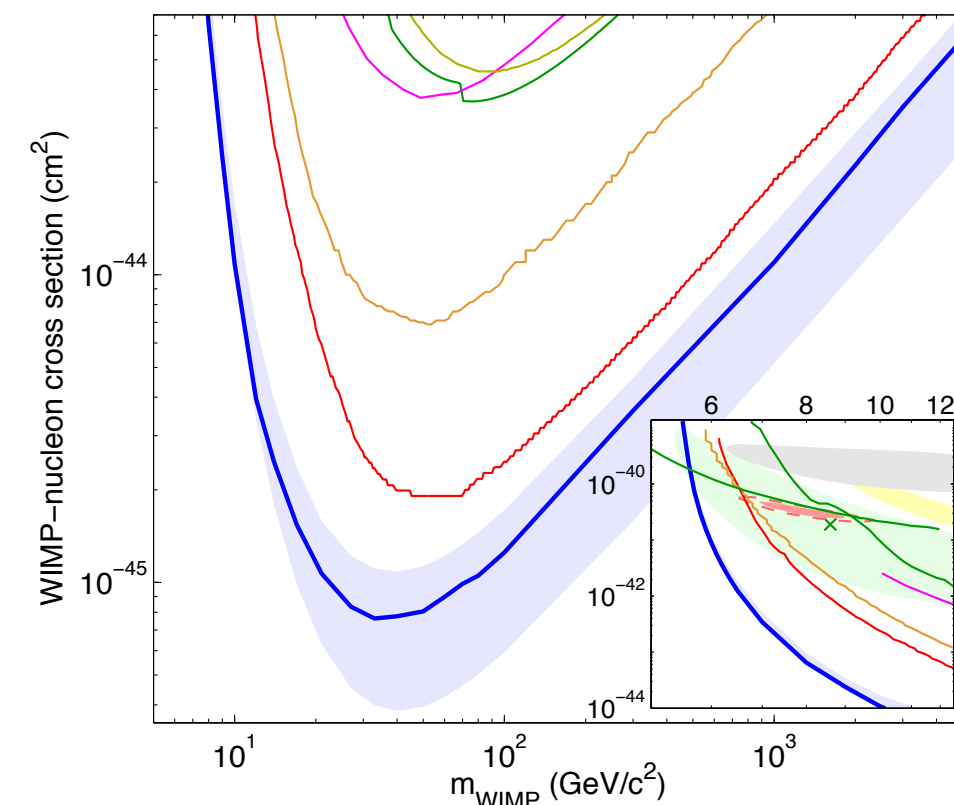
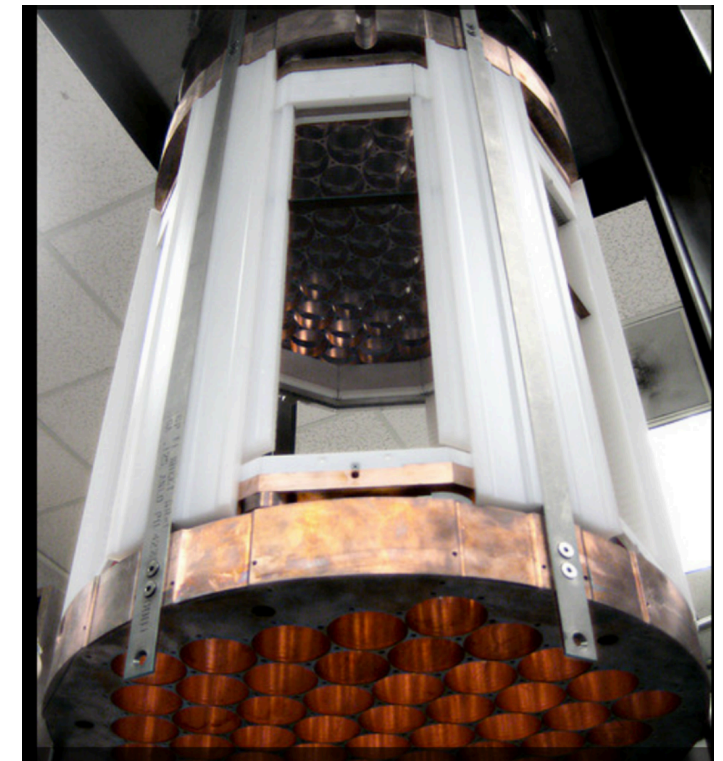
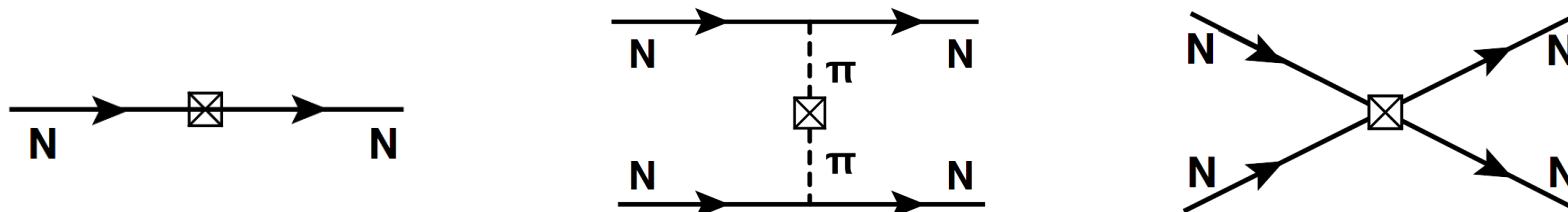
- Dark matter direct detection experiments look for DM interactions with nuclei (Si, Xe, ...)
- One possible interaction is through scalar exchange

$$\mathcal{L} = \frac{G_F}{2} \sum_q a_S^{(q)} (\bar{\chi} \chi) (\bar{q} q)$$

- Accessible via Feynman-Hellman theorem
- At hadronic/nuclear level

$$\mathcal{L} \rightarrow G_F \bar{\chi} \chi \left( \frac{1}{4} \langle 0 | \bar{q} q | 0 \rangle \text{Tr} [a_S \Sigma^\dagger + a_S^\dagger \Sigma] + \frac{1}{4} \langle N | \bar{q} q | N \rangle N^\dagger N \text{Tr} [a_S \Sigma^\dagger + a_S^\dagger \Sigma] - \frac{1}{4} \langle N | \bar{q} \tau^3 q | N \rangle (N^\dagger N \text{Tr} [a_S \Sigma^\dagger + a_S^\dagger \Sigma] - 4 N^\dagger a_{S,\xi} N) + \dots \right)$$

- Contributions:

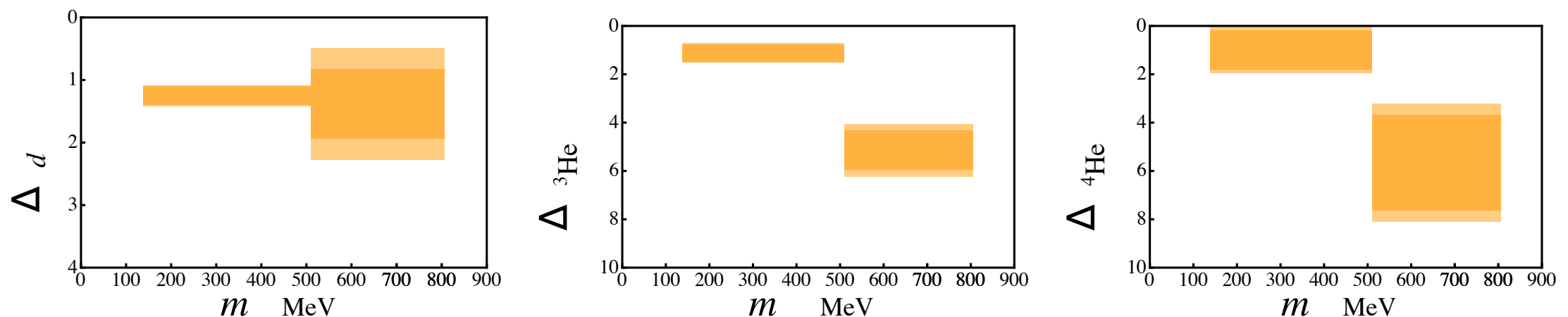


# Nuclear sigma terms

- Previous work suggested scalar-isoscalar dark matter couplings to nuclei have  $O(50\%)$  uncertainty arising from MECs [Prezeau et al 2003]
- Quark mass dependence of nuclear binding energies bounds such contributions

$$\delta\sigma_{Z,N} = \frac{\langle Z, N(\text{gs}) | \bar{u}u + \bar{d}d | Z, N(\text{gs}) \rangle}{A \langle N | \bar{u}u + \bar{d}d | N \rangle} - 1 = -\frac{1}{A\sigma_N} \frac{m_\pi}{2} \frac{d}{dm_\pi} B_{Z,N}$$

- Lattice calculations + physical point suggest such contributions are  $O(10\%)$  or less for light nuclei

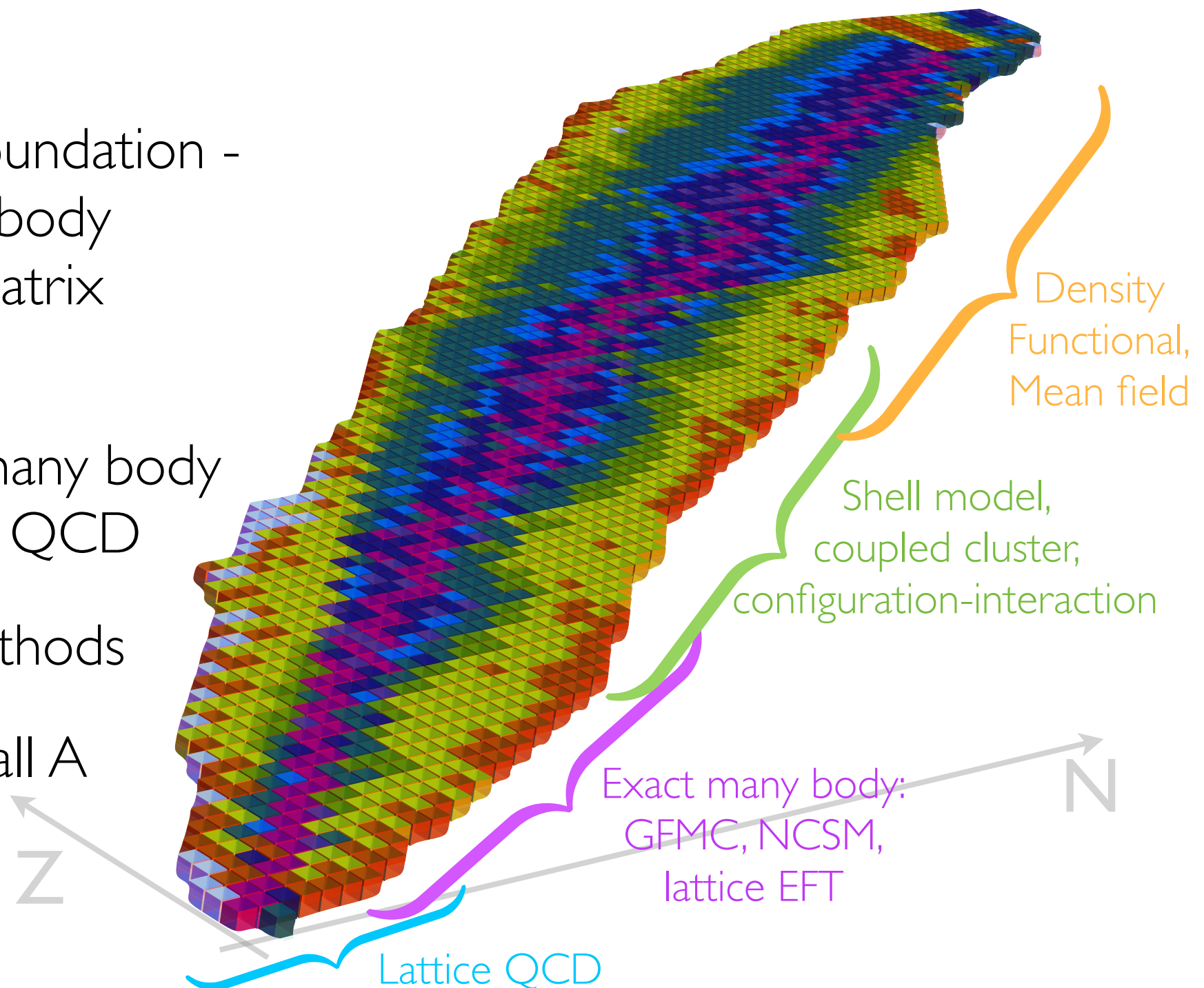


- Admittedly crude approximation to derivative ... stay tuned



# Larger nuclei

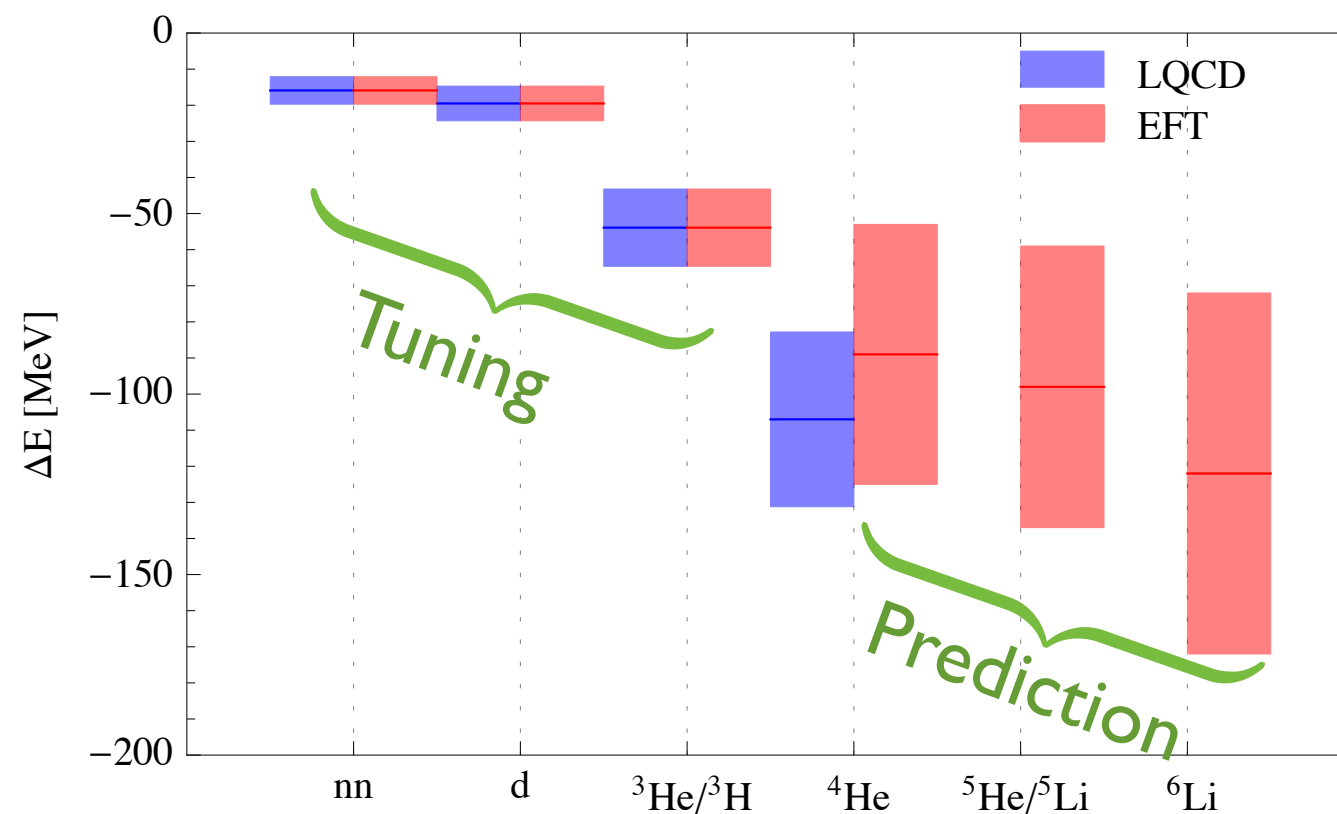
- A path to ab initio nuclear physics:
- QCD forms a foundation - determines few body interactions & matrix elements
- Match existing many body techniques onto QCD
- Hierarchy of methods
- QCD: focus on small  $A$
- ... for now ...



# Heavy quark universe

[Barnea, et al. 1311.4966]

- Already seeing LQCD and nuclear EFT coming together
- For heavy quarks, even spectroscopy requires QCD matching



- Equally important for matrix elements at the physical quark mass

# *Matrix elements: philosophy*

- Provided the hierarchy of higher-body interactions persists into heavy nuclei, power counting of nuclear effective field theory:
  - 1-body currents are dominant
  - 2-body currents are sub-leading and higher-body currents are even less important
- Determine one body contributions from single nucleon/pion systems
- Determine few-body contributions from  $A=2,3,4\dots$
- Use EFT and many body methods to extend LQCD results to large nuclei



# *Prospects and Outlook*

- Properties and interactions of light nuclei represent an important opportunity for LQCD
- Direct impact on NP of light nuclei
- Input to/constraint on nuclear many-body methods can greatly improve NP of searches for BSM physics
- Stay tuned



[FIN]

# NuTeV anomaly: $\sin^2\theta_w$

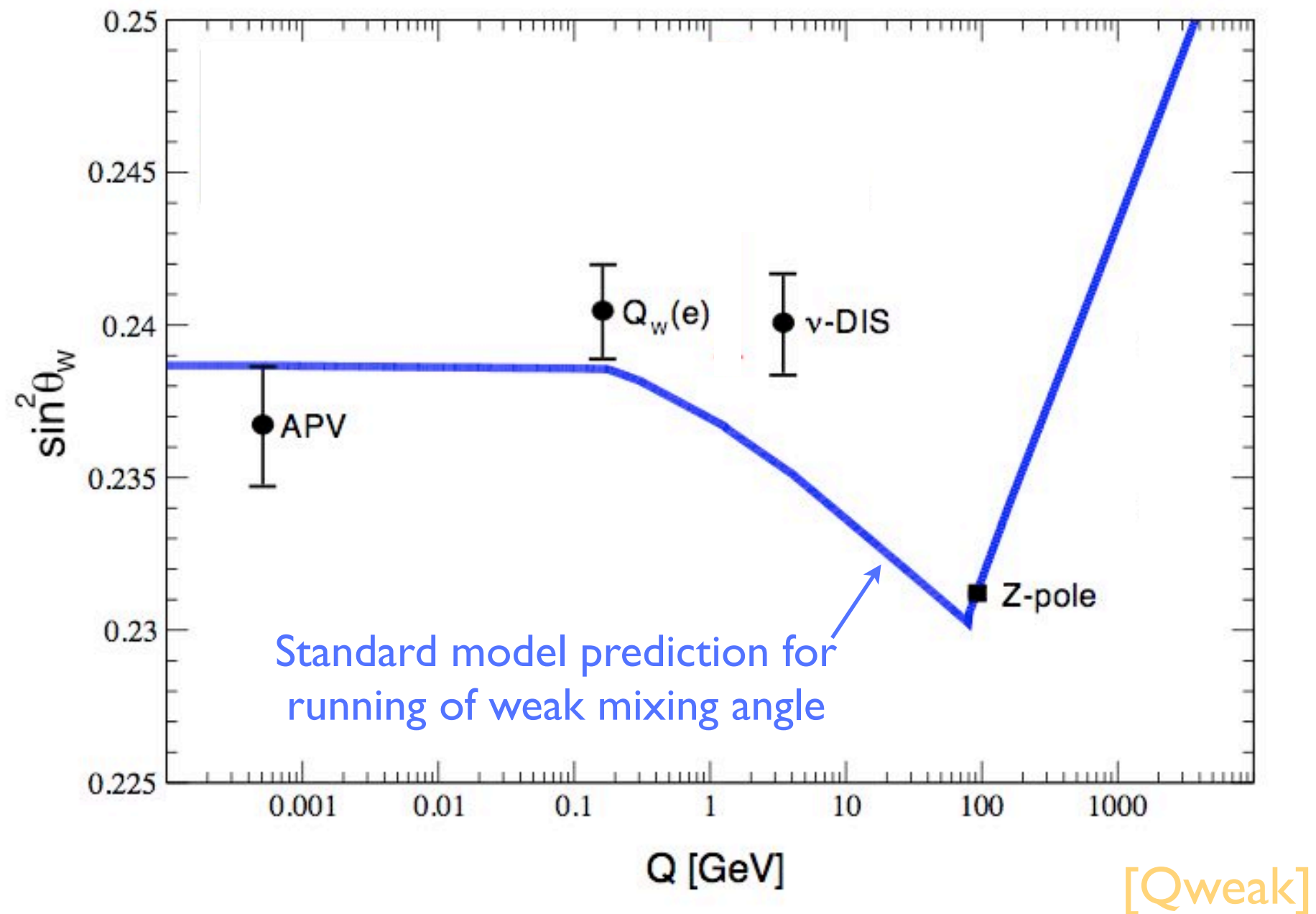
- Neutrino deep-inelastic scattering

- Paschos-Wolfenstein relation:

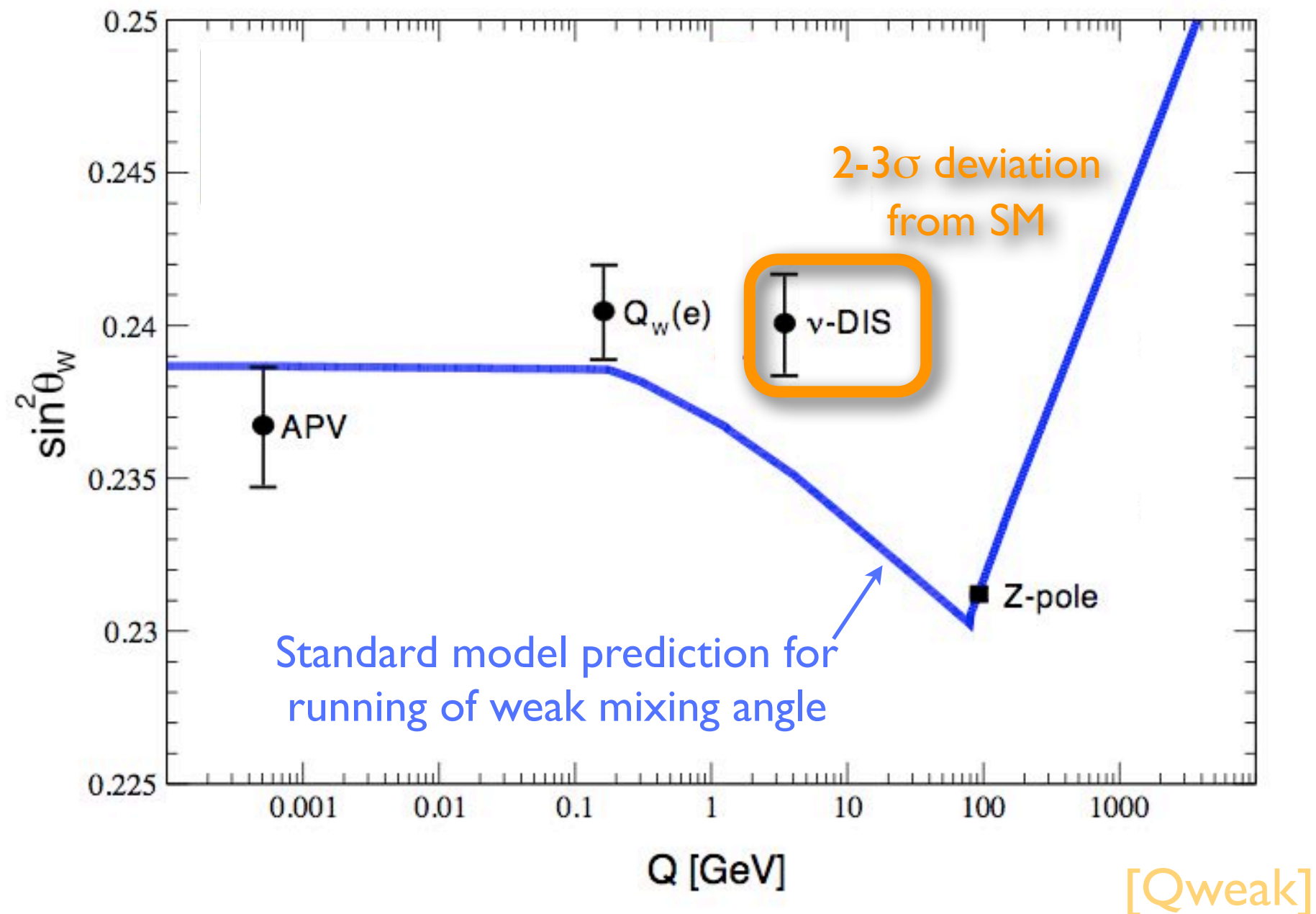
$$R^- = \frac{\sigma_{NC}^{\nu N} - \sigma_{NC}^{\bar{\nu} N}}{\sigma_{CC}^{\nu N} - \sigma_{CC}^{\bar{\nu} N}} = \frac{1}{2} - \sin^2 \theta_w$$

- NuTeV measure CC and NC neutrino scattering on steel target at Fermilab
- Extract the weak mixing angle

# *NuTeV anomaly: $\sin^2\theta_w$*



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# NuTeV anomaly: $\sin^2 \theta_W$

- Corrected Paschos-Wolfenstein relation:

$$R_A^- = \frac{\sigma_{NC}^{\nu A} - \sigma_{NC}^{\bar{\nu} A}}{\sigma_{CC}^{\nu A} - \sigma_{CC}^{\bar{\nu} A}} = \frac{1}{2} - \sin^2 \theta_W + \epsilon_v + \epsilon_n + \epsilon_s + \epsilon_c$$

Nuclear modifications  
Non-isoscalarity

Strange quarks  
Charm quarks

- NuTeV take some of this into account
- Many authors find significant reduction in NuTeV significance from hadronic/nuclear physics

# EMC for light nuclei

- E03-103 experiment @ JLab
- High precision studies of DIS on light nuclei
- Size of EMC effect vs nuclear density
- Nontrivial behaviour
- Correlated with strength of short range correlations

[Weinstein *et al*, PRL 106 (2011) 052301]

